

CFD - Lecture 4

- Conservation of momentum
- Burger's equation
- An introduction to compressible Euler equations
- Foam Tutorial

- First, everybody log in to CCV.

Conservation of Momentum

$$\frac{d}{dt} (\text{momentum in system}) = \sum (\text{Forces acting on system})$$

- think $F = ma$ from high school

Let $\rho = \text{density}$ [m/L^3]

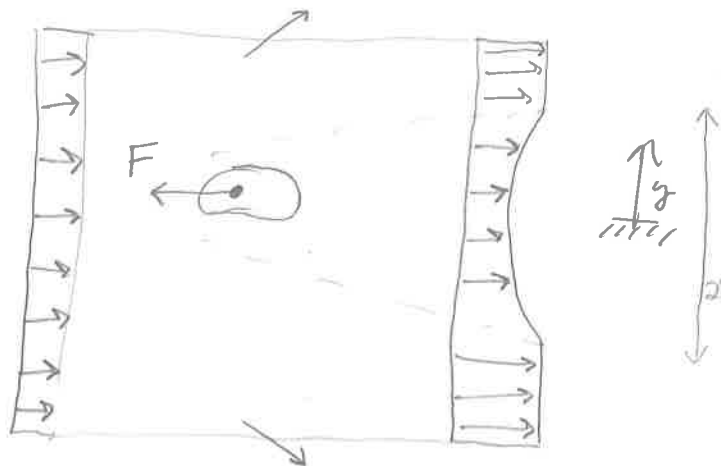
$u = \text{velocity}$ [L/T]

$$\frac{\partial}{\partial t} \int_{\Omega} \rho u \, dx + \oint_{\partial\Omega} (\rho u) u \cdot \hat{n} \, dA = \sum_i F_i$$

\uparrow
non-linear

Ex Drag on a cylinder in a wind tunnel

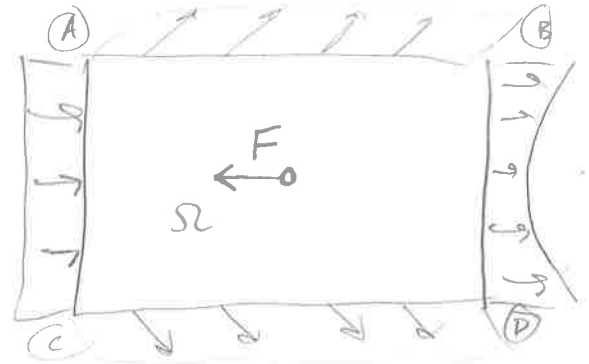
- Assume uniform velocity at inlet U_{∞}
- symmetric velocity at top & bottom
- measured velocity profile at outlet $\rightarrow u(y)$
- Because of drag acting on obstacle, velocity at outlet is slower, causing mass flux at top and bottom
- By symmetry, assume equal & opposite



- From conservation, can recover force on obstacle from outlet vel.

• assume $u(y) = \begin{cases} U_{\infty} & |y| > h \\ U_{\infty} \left(1 - \frac{1}{2} \left(\frac{y}{h}\right)^2\right) & |y| < h \end{cases}$

$$\cancel{\int_{\Omega} \rho u dx} + \int_{\partial\Omega} (\rho u) \cdot (u \cdot \hat{n}) dA = -F$$



$$\int_{AC} (\rho u) u \cdot \hat{n} dA + \int_{BD} (\rho u) u \cdot \hat{n} dA = -F$$

$$\rho \left(\int_{-h}^h U_{\infty} (-U_{\infty}) dy + \int_{-h}^h u(y)^2 dy \right) = -F$$

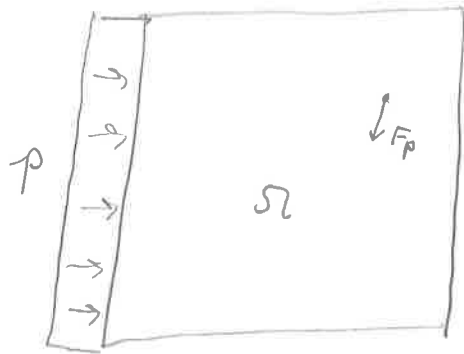
$$\rho \left(-U_{\infty}^2 \int_{-h}^h dy + \int_{-h}^h \left[U_{\infty} \left(1 - \frac{1}{2} \left(\frac{y}{h}\right)^2\right) \right]^2 dy \right) = -F$$

$$\rho \left(-2h U_{\infty}^2 + \frac{43}{30} h U_{\infty}^2 \right) = -F$$

$$F = \frac{17}{30} \rho h U_{\infty}^2$$

In general, forces can take form of:
 point force, surface force, body force

$$F = F_p + \int_{\partial\Omega} p \, dA + \int_{\Omega} f_B \, dV$$



We'll see more about this next week
 when we discuss Navier-Stokes derivation

Burgers Equation

- Taking $\sum F = 0$, $\ell = 1$

gives Burgers Egn
 up to factor of 1/2

$$\partial_t u + \partial_x \left(\frac{1}{2} u^2 \right) = 0$$

- Burgers equation is a model for 1D non-linear advection \rightarrow exact solutions mean that its easy to prove/demonstrate convergence

- Expand flux term to see

$$\partial_t u + u \partial_x u = 0$$

compare to advection eqn \rightarrow

$$\partial_t u + a \partial_x u = 0$$

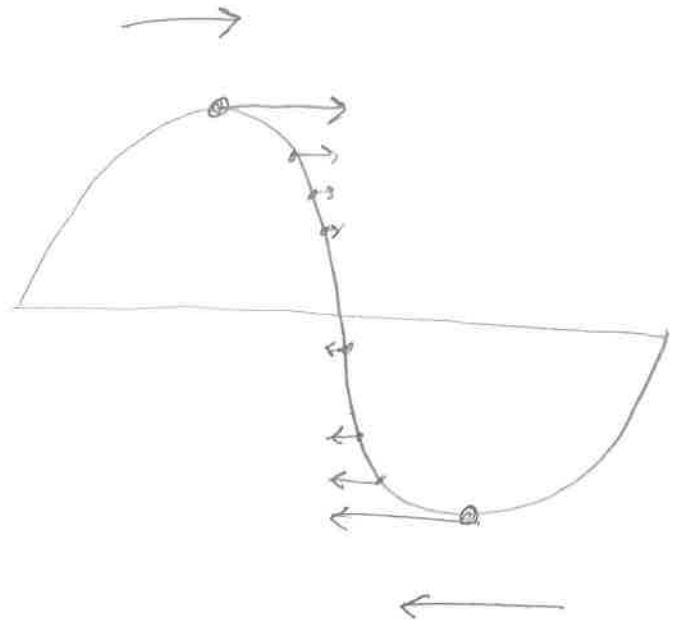
- Non-linear advection

- 1D model problem

- Take $u = \sin x$

- Advection magnitude is non-uniform and depends on solution

- Solve the exact same way, but we'll just be careful about picking the right upstream value



- HW3

Compressible Euler Equations (CEE)

- What we have now is a suite of techniques for handling a variety of physics

$$\text{CEE} \left\{ \begin{array}{l} \partial_t(\rho u) + \nabla \cdot (\rho u^2) = -\nabla p + \mu \nabla^2 u \\ \partial_t \rho + \nabla \cdot (\rho u) = 0 \\ \rho = \rho(p, T) \quad \text{ex} \quad \rho = \frac{p}{RT} \end{array} \right.$$

Next week we derive these, but first we'll warm up in Foam

Foam tutorial

- In last class we did flow through a pipe.
- For easy geometry, I'll include calculation for flow between infinite plates here

$$u_{in} = u_{\infty}$$

$$u_{out} = C y \left(1 - \frac{y}{L}\right)$$

$$\int_{\Omega} \vec{u} \cdot \hat{n} dA = 0$$

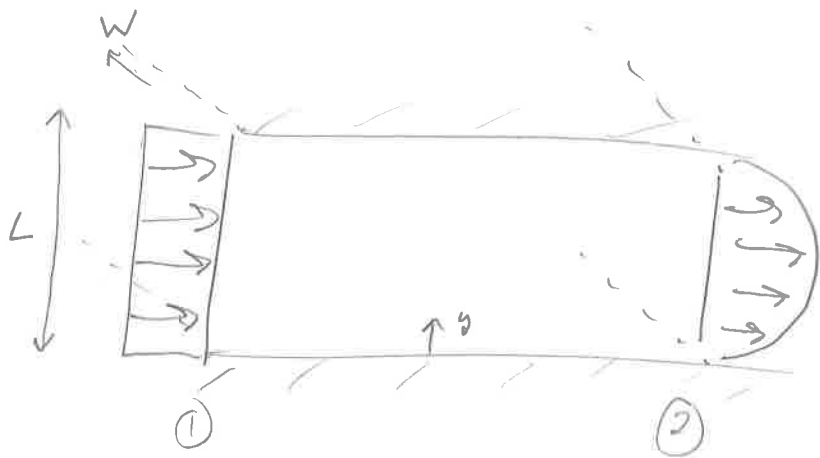
$$\int_{(1)} u \cdot \hat{n} dA + \int_{(2)} u \cdot \hat{n} dA = 0$$

$$dA = W dy$$

$$- \int_0^L u_{\infty} W dy + \int_0^L C y \left(1 - \frac{y}{L}\right) W dy = 0$$

$$- u_{\infty} L + C \int_0^L y \left(1 - \frac{y}{L}\right) dy = 0$$

$$- u_{\infty} L + \frac{C L^2}{6} = 0 \rightarrow C = \frac{6 u_{\infty}}{L}$$



$$u_{out}(y) = \frac{6 u_{\infty}}{L^2} y(L-y)$$