

A SUPERCONVERGENT LDG-HYBRIDIZABLE GALERKIN METHOD FOR SECOND-ORDER ELLIPTIC PROBLEMS

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ABSTRACT. We identify and study an LDG-hybridizable Galerkin method, which is not an LDG method, for second-order elliptic problems in several space dimensions with remarkable convergence properties. Unlike all other known discontinuous Galerkin methods using polynomials of degree $k \geq 0$ for both the potential as well as the flux, the order of convergence in L^2 of both unknowns is $k + 1$. Moreover, both the approximate potential as well as its numerical trace superconverge in L^2 -like norms, to suitably chosen projections of the potential, with order $k + 2$. This allows the application of element-by-element postprocessing of the approximate solution which provides an approximation of the potential converging with order $k + 2$ in L^2 . The method can be thought to be *in between* the hybridized version of the Raviart-Thomas and that of the Brezzi-Douglas-Marini mixed methods.

1. INTRODUCTION

In this paper, we consider the LDG-hybridizable (LDG-H) Galerkin methods recently introduced in [9] and show how to define their numerical traces in order to achieve the optimal order of convergence for the approximation to the flux, and to obtain superconvergence properties similar to those of the hybridized mixed methods of Raviart-Thomas (RT) [13] and the Brezzi-Douglas-Marini (BDM) [4] methods; see also [8].

For the sake of simplicity of the exposition, we carry out this in the setting of the model second-order elliptic problem

$$(1.1a) \quad \mathbf{c} \mathbf{q} + \nabla u = 0 \quad \text{in } \Omega,$$

$$(1.1b) \quad \nabla \cdot \mathbf{q} = f \quad \text{in } \Omega,$$

$$(1.1c) \quad u = g \quad \text{on } \partial\Omega_D,$$

$$(1.1d) \quad \mathbf{q} \cdot \mathbf{n} = \mathbf{q}_N \quad \text{on } \partial\Omega_N,$$

where $\Omega \subset \mathbb{R}^d$ is a polyhedral domain ($d \geq 2$), $f \in L^2(\Omega)$, and $\mathbf{c} = \mathbf{c}(\mathbf{x})$ is a symmetric $d \times d$ matrix function that is uniformly positive definite on Ω with components in $L^\infty(\Omega)$. As usual, we assume that the $(d - 1)$ -Lebesgue measure of $\partial\Omega_D$ is not zero, that $\partial\Omega = \overline{\partial\Omega_D} \cup \overline{\partial\Omega_N}$ and that $\partial\Omega_D \cap \partial\Omega_N = \emptyset$.

To describe our results, we need to introduce the hybridized Galerkin methods studied in [9]. To do that, let us introduce some notation. We denote by $\Omega_h = \{K\}$

1991 *Mathematics Subject Classification.* 65M60,65N30,35L65.

Key words and phrases. discontinuous Galerkin methods, hybridization, superconvergence, second-order elliptic problems.

The first author was supported in part by the National Science Foundation (Grant DMS-0411254) and by the University of Minnesota Supercomputing Institute. The third author was supported by an NSF Mathematical Science Postdoctoral Research Fellowship (DMS-0503050).

a triangulation of the domain Ω of shape-regular simplexes K and set $\partial\Omega_h := \{\partial K : K \in \Omega_h\}$. We associate to this triangulation the set of interior faces \mathcal{E}_h^i and the set of boundary faces \mathcal{E}_h^∂ . We say that $e \in \mathcal{E}_h^i$ if there are two simplexes K^+ and K^- in Ω_h such that $e = \partial K^+ \cap \partial K^-$, and we say that $e \in \mathcal{E}_h^\partial$ if there is a simplex in Ω_h such that $e = \partial K \cap \partial\Omega$. We set $\mathcal{E}_h := \mathcal{E}_h^i \cup \mathcal{E}_h^\partial$.

The hybridized Galerkin methods seek an approximation to the exact solution $(\mathbf{q}|_\Omega, u|_\Omega, u|_{\mathcal{E}_h \setminus \partial\Omega_N})$, $(\mathbf{q}_h, u_h, \lambda_h)$, in a finite-dimensional space $\mathbf{V}_h \times W_h \times M_h$ of the form

$$(1.2a) \quad \mathbf{V}_h := \{\mathbf{v} \in \mathbf{L}^2(\Omega) : \mathbf{v}|_K \in \mathbf{V}(K) \quad \forall K \in \Omega_h\},$$

$$(1.2b) \quad W_h := \{\omega \in L^2(\Omega) : \omega|_K \in W(K) \quad \forall K \in \Omega_h\},$$

$$(1.2c) \quad M_h := \{\mathbf{m} \in L^2(\partial\Omega_h) : \mathbf{m}|_e \in M(e) \quad \forall e \in \mathcal{E}_h, \quad \mathbf{m}|_{\partial\Omega_D} = 0\},$$

and determines it by requiring that

$$(1.3a) \quad (\mathbf{c} \mathbf{q}_h, \mathbf{v})_{\Omega_h} - (u_h, \nabla \cdot \mathbf{v})_{\Omega_h} + \langle \widehat{u}_h, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial\Omega_h} = 0,$$

$$(1.3b) \quad -(\mathbf{q}_h, \nabla \omega)_{\Omega_h} + \langle \widehat{\mathbf{q}}_h \cdot \mathbf{n}, \omega \rangle_{\partial\Omega_h} = (f, \omega)_{\Omega_h},$$

$$(1.3c) \quad \langle \widehat{\mathbf{q}}_h \cdot \mathbf{n}, \mu \rangle_{\partial\Omega_h} = \langle \mathbf{q}_N, \mu \rangle_{\partial\Omega_N},$$

for all $(\mathbf{v}, \omega, \mu) \in \mathbf{V}_h \times W_h \times M_h$. Here, we have used the notation

$$\begin{aligned} (\boldsymbol{\sigma}, \mathbf{v})_{\Omega_h} &:= \sum_{K \in \Omega_h} \int_K \boldsymbol{\sigma}(x) \cdot \mathbf{v}(x) \, dx, \\ (\zeta, \omega)_{\Omega_h} &:= \sum_{K \in \Omega_h} \int_K \zeta(x) \omega(x) \, dx, \\ \langle \zeta, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial\Omega_h} &:= \sum_{K \in \Omega_h} \int_{\partial K} \zeta(\gamma) \mathbf{v}(\gamma) \cdot \mathbf{n} \, d\gamma, \end{aligned}$$

for any functions $\boldsymbol{\sigma}, \mathbf{v}$ in $\mathbf{H}^1(\Omega_h) := [H^1(\Omega_h)]^d$ and ζ, ω in the space $H^1(\Omega_h) = \{v \in L^2(\Omega) : v|_K \in H^1(K) \quad \forall K \in \Omega_h\}$. The outward normal unit vector to ∂K is denoted by \mathbf{n} .

To complete the description of the hybridized Galerkin methods, the definition of numerical traces $(\widehat{\mathbf{q}}_h, \widehat{u}_h)$ on the faces of the triangulation \mathcal{E}_h has to be provided. The choice which is relevant here is

$$(1.4a) \quad \widehat{u}_h = \begin{cases} \mathbf{P}_\partial g & \text{on } \mathcal{E}_h \cap \partial\Omega_D, \\ \lambda_h & \text{on } \mathcal{E}_h \setminus \partial\Omega_D, \end{cases}$$

$$(1.4b) \quad \widehat{\mathbf{q}}_h = \mathbf{q}_h + \tau (u_h - \widehat{u}_h) \mathbf{n} \quad \text{on } \mathcal{E}_h,$$

where \mathbf{P}_∂ denotes an L^2 -projection defined as follows. Given any function $\zeta \in L^2(\mathcal{E}_h)$ and an arbitrary face $e \in \mathcal{E}_h$, the restriction of $\mathbf{P}_\partial \zeta$ to e is defined as the element of $\mathcal{P}^k(e)$ that satisfies

$$(1.5) \quad \langle \mathbf{P}_\partial \zeta - \zeta, \omega \rangle_e = 0, \quad \forall \omega \in \mathcal{P}^k(e).$$

Note that by suitably choosing the *local spaces* $\mathbf{V}(K)$, $W(K)$, and $M(e)$, and the values of the *local stabilization parameters* τ , we can obtain the hybridized RT $_k$, the hybridized BDM $_k$ and the LDG-H $_k$ methods; see Tables 1 and 2. In Table 1 and in the remaining of this paper, we denote the space of polynomials of degree at most $k \geq 0$ defined on D by $\mathcal{P}^k(D)$, and set $\mathcal{P}^k(D) := [\mathcal{P}^k(D)]^d$. Since all these methods can be implemented in the same way and can be used in different

elements while being automatically coupled, what is relevant, as argued in [9], is to find out which method should be used in what *element* in order to optimize the computational effort. It is thus important from this perspective to develop DG methods as accurate and efficient as mixed methods so that they could be used in situations in which mixed method cannot. The LDG-H methods we uncover in this paper are the first example of those methods.

TABLE 1. The local spaces

method	$\mathbf{V}(K)$	$W(K)$	$M(e)$
RT $_k$	$\mathcal{P}^k(K) \oplus \mathbf{x} \mathcal{P}^k(K)$	$\mathcal{P}^k(K)$	$\mathcal{P}^k(e)$
LDG-H $_k$	$\mathcal{P}^k(K)$	$\mathcal{P}^k(K)$	$\mathcal{P}^k(e)$
BDM $_k$	$\mathcal{P}^k(K)$	$\mathcal{P}^{k-1}(K)$	$\mathcal{P}^k(e)$

TABLE 2. The local stabilization parameters τ

method	$\tau _{\partial K}$
RT $_k$	$\equiv 0$
LDG-H $_k$	$\geq 0, \neq 0$
BDM $_k$	$\equiv 0$

It is well known that the RT $_k$ and BDM $_k$ methods provide an approximation \mathbf{q}_h to the flux which converges in L^2 with order $k+1$, that u_h and λ_h superconverge in L^2 -like norms to suitably chosen projections of the potential u with order $k+2$, and that, as a consequence, it is possible to postprocess the approximate solution to obtain another approximation u_h^* converging in L^2 with order $k+2$; see [1] and [4], and also [8]. In this paper, we use an extension of the postprocessing proposed in [14, 15] and [12]. Given the similarities between these two mixed methods and the LDG-H $_k$ method, it is natural to ask if it is possible to choose the local penalization parameters τ as to obtain similar convergence and superconvergence results. The main contribution of this paper is to show that this is actually possible.

Indeed, we show that this happens if we take, on each simplex $K \in \Omega_h$

$$(1.6) \quad \tau = \begin{cases} 0, & \text{on } \partial K \setminus e_K^\tau, \\ \tau_K, & \text{on } e_K^\tau, \end{cases}$$

where e_K^τ is an arbitrary but fixed face of K and τ_K is a strictly positive real number. Due to the superconvergence properties of this DG method, we say that this is an SCDG method; for simplicity, we are going to refer to the method under consideration by the SCDG $_k$ method. It is interesting to note two of the *minimal dissipation* DG methods considered in [6], in the framework of a study of superconvergence properties of DG methods for one-dimensional steady-state convection-diffusion problems, happen to be an SCDG method. The first is called

the md-DG method, see Table 1 in [6], and is obtained, in our notation, by taking on each interior node x_i ,

$$\tau(x_i^+) = 0 \quad \text{and} \quad \tau(x_i^-) = k/h_i,$$

where h_i is the size of the interval to the left of the node x_i . The second is called the md-LDG method, and is obtained by taking the above choice of parameters τ formally letting $\tau(x_i^-)$ go to infinity. The authors are not aware of any other instance of SCDG methods. In particular, let us emphasize that SCDG *methods are not* LDG *methods* whenever the stabilization parameters τ are finite; see the discussion about LDG-H methods in [9].

In Table 3, we compare the orders of convergence for the flux of this method and the above-mentioned mixed methods. We have also included the order of convergence for the general LDG-H methods; it can be deduced from their characterization [9] and the study of DG methods carried out in [5]. Finally, in Table 4, we display the orders of convergence of the postprocessed approximation u_h^* to the potential.

TABLE 3. The orders of convergence in h of the L^2 -errors

method	$\ \mathbf{q} - \mathbf{q}_h\ _{L^2(\Omega_h)}$	$\ u - u_h\ _{L^2(\Omega_h)}$	condition
RT $_k$	$k + 1$	$k + 1$	$k \geq 0$
LDG-H $_k$	k	$k + 1$	$k \geq 1$ and $\tau = \mathcal{O}(1/h)$
LDG-H $_k$	$k + 1/2$	$k + 1$	$k \geq 0$ and $\tau = \mathcal{O}(1)$
SCDG $_k$	$k + 1$	$k + 1$	$k \geq 0$
BDM $_k$	$k + 1$	k	$k \geq 1$

TABLE 4. The orders of convergence in h of $\|u - u_h^*\|_{L^2(\Omega_h)}$

method	order	condition
RT $_k$	$k + 2$	$k \geq 0$
SCDG $_k$	$k + 2$	$k \geq 1$
SCDG $_k$	$k + 2$	$k = 0, f = 0$
BDM $_k$	$k + 2$	$k \geq 2$
BDM $_k$	$k + 1$	$k = 1$

We also uncover new relations between these three methods. One of the main features of the hybridized Galerkin methods proposed in [9] is that the only degrees of freedom that turn out to be globally coupled are those of the so-called Lagrange multiplier λ_h . This implies, in particular, that the LDG-H methods can be *more efficiently* implemented than the LDG methods introduced in [10]. In fact, they can be implemented as efficiently as the hybridized RT $_k$ and BDM $_k$ mixed methods; see [9] for details. Here we show that the stiffness matrix of the Lagrange multiplier for the RT $_k$, BDM $_k$ and SCDG $_k$ methods is actually *identical* and that, when $f|_K \in \mathcal{P}^{k-1}(K)$ for all $K \in \Omega_h$, these methods provide the *same* approximation $(\mathbf{q}_h, \lambda_h)$.

Next, let us briefly comment on the approach taken to carry out the a priori error analysis of the SCDG $_k$ methods. We did not take the approach used in [5] to analyze DG and LDG methods, or that used in the unified analysis of DG methods [2]. Instead, we exploited the unifying framework of hybridized Galerkin methods [9] to render the analysis of the SCDG $_k$ methods as close as possible to those of the hybridized RT $_k$ and BDM $_k$ methods. Since a key ingredient in those analyzes is the existence of a projection $(\mathbf{\Pi}, \mathbb{P})$ satisfying the so-called *commutativity* property

$$\nabla \cdot \mathbf{\Pi}\boldsymbol{\sigma} = \mathbb{P}^t \nabla \cdot \boldsymbol{\sigma},$$

for all $\boldsymbol{\sigma} \in \mathbf{H}(\text{div}, \Omega)$, the crucial step in the analysis was to find a similar projection. Unlike the above-mentioned mixed methods, the space of fluxes \mathbf{V}_h of the SCDG $_k$ methods is not included in $\mathbf{H}(\text{div}, \Omega)$ and, as a consequence, the above commutativity property can only be satisfied in a weak sense. We found a *new* projection satisfying the following weak version of the commutativity property:

$$-(\nabla\zeta, \mathbf{\Pi}\boldsymbol{\sigma})_{\Omega_h} = (\mathbb{P}\zeta, \nabla \cdot \boldsymbol{\sigma})_{\Omega_h}$$

for all $(\boldsymbol{\sigma}, \zeta) \in \mathbf{H}^1(\Omega_h) \times H^1(\Omega_h)$ such that $\zeta|_{\partial\Omega} = 0$. Just as the local spaces of the SCDG $_k$ methods are, roughly speaking, “in between” the local spaces of the RT $_k$ and BDM $_k$ methods, this projection can also be considered to be “in between” the corresponding projections of those mixed methods. The construction of this projection, which is intimately linked to the definition of the numerical traces of the method \widehat{u}_h and $\widehat{\mathbf{q}}_h$ and to the choice of the local spaces, is certainly the most interesting aspect of the analysis of the SCDG $_k$ methods. The first component of the projection, $\mathbf{\Pi}$, was used in the error analysis of the *minimal dissipation* LDG method in [7].

The paper is organized as follows. In Section 2, we state and discuss our main results and then prove them in Section 3. In Section 4, we display numerical experiments validating the theoretical results. Finally, in Section 5, we end with some concluding remarks.

2. THE MAIN RESULTS

2.1. The projection $(\mathbf{\Pi}, \mathbb{P})$. In this subsection, we define the projection

$$(\mathbf{\Pi}, \mathbb{P}) : \mathbf{H}^1(\Omega_h) \times H^1(\Omega_h) \rightarrow \mathbf{V}_h \times W_h,$$

and gather its main properties.

Given a function $\boldsymbol{\sigma} \in \mathbf{H}^1(\Omega_h)$ and an arbitrary simplex $K \in \Omega_h$, the restriction of $\mathbf{\Pi}\boldsymbol{\sigma}$ to K is defined as the element of $\mathcal{P}^k(K)$ that satisfies

$$(2.7a) \quad (\mathbf{\Pi}\boldsymbol{\sigma} - \boldsymbol{\sigma}, \mathbf{v})_K = 0, \quad \forall \mathbf{v} \in \mathcal{P}^{k-1}(K), \text{ if } k \geq 1,$$

$$(2.7b) \quad \langle (\mathbf{\Pi}\boldsymbol{\sigma} - \boldsymbol{\sigma}) \cdot \mathbf{n}, \omega \rangle_e = 0, \quad \forall \omega \in \mathcal{P}^k(e) \text{ and } \forall e \in \partial K, e \neq e_K^\tau.$$

Similarly, given a function $\zeta \in H^1(\Omega_h)$ and an arbitrary simplex $K \in \Omega_h$, the restriction of $\mathbb{P}\zeta$ to K is defined as the element of $\mathcal{P}^k(K)$ that satisfies

$$(2.8a) \quad (\mathbb{P}\zeta - \zeta, \mathbf{w})_K = 0, \quad \forall \mathbf{w} \in \mathcal{P}^{k-1}(K), \text{ if } k \geq 1,$$

$$(2.8b) \quad \langle \mathbb{P}\zeta - \zeta, \omega \rangle_{e_K^\tau} = 0, \quad \forall \omega \in \mathcal{P}^k(e_K^\tau).$$

We gather the main properties of this projection in the following result. To state it, we need to recall the definition of some classical projections. Given a

function $\boldsymbol{\sigma} \in \mathbf{H}^1(\Omega_h)$ and an arbitrary simplex $K \in \Omega_h$, the restriction of $\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma}$ to K is defined as the element of $\mathcal{P}^k(K) \oplus \mathbf{x}\mathcal{P}(K)^k$ that satisfies

$$(2.9a) \quad (\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma} - \boldsymbol{\sigma}, \mathbf{v})_K = 0, \quad \forall \mathbf{v} \in \mathcal{P}^{k-1}(K), \text{ if } k \geq 1,$$

$$(2.9b) \quad \langle (\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma} - \boldsymbol{\sigma}) \cdot \mathbf{n}, \omega \rangle_e = 0, \quad \forall \omega \in \mathcal{P}^k(e), \text{ for all faces of } K.$$

Given a function $\zeta \in H^1(\Omega_h)$ and an arbitrary simplex $K \in \Omega_h$, the restriction of $\mathbf{P}^\ell\zeta$ to K is defined as the element of $\mathcal{P}^\ell(K)$ that satisfies

$$(2.10) \quad (\mathbf{P}^\ell\zeta - \zeta, \omega)_K = 0, \quad \forall \omega \in \mathcal{P}^\ell(K).$$

To simplify the notation, we are going to write \mathbf{P} instead of \mathbf{P}^k . Note that $(\mathbf{\Pi}^{\text{RT}}, \mathbf{P})$ is nothing but the projection for the RT $_k$ method. We are now ready to state our result.

Proposition 2.1. *The projection $(\mathbf{\Pi}, \mathbf{P})$ given by (2.7) and (2.8) is well defined. Moreover, on each simplex $K \in \Omega_h$, it satisfies the orthogonality properties*

$$(i) \quad (\zeta - \mathbf{P}\zeta, \nabla \cdot \mathbf{v})_K = 0,$$

$$(ii) \quad (\boldsymbol{\sigma} - \mathbf{\Pi}\boldsymbol{\sigma}, \nabla\omega)_K = 0,$$

$$(iii) \quad \langle \mathbf{P}\zeta - \mathbf{P}_\partial\zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_e = 0 \quad \text{for all faces } e \text{ of } K,$$

for all $(\mathbf{v}, \omega) \in \mathcal{P}^k(K) \times \mathcal{P}^k(K)$, and the weak commutativity property

$$(iv) \quad -(\nabla\zeta, \mathbf{\Pi}\boldsymbol{\sigma})_K = (\mathbf{P}\zeta, \nabla \cdot \boldsymbol{\sigma})_K - \langle \mathbf{P}_\partial\zeta, \mathbf{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K},$$

for all $(\boldsymbol{\sigma}, \zeta) \in \mathbf{H}^1(\Omega_h) \times H^1(\Omega_h)$. Finally, we have the following approximation estimates

$$(v) \quad \|\mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n}\|_{L^2(e_K^\tau)} \leq Ch_K^{r+1/2} |\mathbf{P}\nabla \cdot \boldsymbol{\sigma}|_{H^r(K)},$$

$$(vi) \quad \|\mathbf{\Pi}\boldsymbol{\sigma} - \mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma}\|_{L^2(K)} \leq Ch_K^{r+1} |\mathbf{P}\nabla \cdot \boldsymbol{\sigma}|_{H^r(K)},$$

$$(vii) \quad \|\mathbf{P}\zeta - \mathbf{P}\zeta\|_{L^2(K)} \leq Ch_K^{s+1} |\nabla\zeta|_{\mathbf{H}^s(K)},$$

where $r, s \in [0, k]$, h_K is the diameter of K , and C depends only on k and the shape-regularity parameters of the simplex K , for any $(\boldsymbol{\sigma}, \zeta) \in \mathbf{H}^1(K) \times H^1(K)$.

We are going to show that the three orthogonality properties imply all the others; they are thus the crucial properties for the analysis. Note also that, by simply adding the identity (iv) over all $K \in \Omega_h$, we obtain the weak commutativity property discussed in the introduction.

2.2. Characterization of the approximate solution. Next we give a characterization of the approximate solution provided by the SCDG $_k$ method. We begin by characterizing the difference between the numerical traces and the traces of the approximate solutions on each simplex.

Proposition 2.2. *For each simplex $K \in \Omega_h$, we have that,*

$$(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} = \tau(u_h - \widehat{u}_h) = \mathbf{P}_\partial\mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} \quad \text{on } \partial K.$$

We see that the jump $(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n}$ is independent of the value of τ whereas the jump $\widehat{u}_h - u_h$ is inversely proportional to τ . Moreover, by the estimate (v) of Proposition 2.1, we have that,

$$\|(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n}\|_{L^2(e_K^\tau)} \leq Ch_K^{r+1/2} |\mathbf{P}f|_{H^r(K)},$$

for any $r \in [0, k]$, and we see that the size of jump under consideration depends solely on the smoothness of $f|_K$. For example, if $\mathbf{P}f|_K$ is a polynomial of degree $k-1$, then $(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} = 0$ on e_K^τ . This implies that $(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} = 0$ on ∂K for every $K \in \Omega_h$ and, as a consequence, that $\mathbf{q}_h \in \mathbf{H}(\text{div}, \Omega)$. Now, if $f \in H^r(K)$, for some $r \in [0, k]$, then we have that

$$\|(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n}\|_{L^2(e_K^\tau)} \leq C h_K^{r+1/2} |f|_{H^r(K)},$$

by well-known approximation properties of the projection \mathbf{P} .

Next, we give a characterization of the approximate solution which follows from a similar result for general hybridized Galerkin methods obtained in [9]. To state it, we need to introduce the *local solvers* associated with the method. The first local solver is defined on the simplex $K \in \Omega_h$ as the mapping $\mathbf{m} \in L^2(\partial K) \rightarrow (\mathbf{Q}\mathbf{m}, \mathcal{U}\mathbf{m}) \in \mathcal{P}^k(K) \times \mathcal{P}^k(K)$ where

$$(2.11a) \quad (\mathbf{c}\mathbf{Q}\mathbf{m}, \mathbf{v})_K - (\mathcal{U}\mathbf{m}, \nabla \cdot \mathbf{v})_K = -\langle \mathbf{m}, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K},$$

$$(2.11b) \quad -(\nabla w, \mathbf{Q}\mathbf{m})_K + \langle w, \widehat{\mathbf{Q}}\mathbf{m} \cdot \mathbf{n} \rangle_{\partial K} = 0,$$

for all $(\mathbf{v}, w) \in \mathcal{P}^k(K) \times \mathcal{P}^k(K)$, where

$$(2.11c) \quad \widehat{\mathbf{Q}}\mathbf{m} = \mathbf{Q}\mathbf{m} + \tau(\mathcal{U}\mathbf{m} - \mathbf{P}_{\partial K}\mathbf{m})\mathbf{n}.$$

The other local solver is defined on the simplex $K \in \Omega_h$ as the mapping $f \in L^2(K) \rightarrow (\mathbf{Q}f, \mathcal{U}f) \in \mathcal{P}^k(K) \times \mathcal{P}^k(K)$ where

$$(2.12a) \quad (\mathbf{c}\mathbf{Q}f, \mathbf{v})_K - (\mathcal{U}f, \nabla \cdot \mathbf{v})_K = 0,$$

$$(2.12b) \quad -(\nabla w, \mathbf{Q}f)_K + \langle w, \widehat{\mathbf{Q}}f \cdot \mathbf{n} \rangle_{\partial K} = (f, w)_K,$$

for all $(\mathbf{v}, w) \in \mathcal{P}^k(K) \times \mathcal{P}^k(K)$, where

$$(2.12c) \quad \widehat{\mathbf{Q}}f = \mathbf{Q}f + \tau \mathcal{U}f \mathbf{n}.$$

We can now state our characterization result.

Theorem 2.3. *The approximate solution $(\mathbf{q}_h, u_h, \lambda_h) \in \mathbf{V}_h \times W_h \times M_h$ given by the SCDG $_k$ method is well defined. Moreover, we have that*

$$(\mathbf{q}_h, u_h) = (\mathbf{Q}\lambda_h, \mathcal{U}\lambda_h) + (\mathbf{Q}g, \mathcal{U}g) + (\mathbf{Q}f, \mathcal{U}f),$$

where λ_h can be characterized as the function in M_h satisfying

$$a_h(\lambda_h, \mu) = b_h(\mu) \quad \forall \mu \in M_h,$$

where

$$\begin{aligned} a_h(\eta, \mu) &:= (\mathbf{c}\mathbf{Q}\eta, \mathbf{Q}\mu)_{\Omega_h}, \\ b_h(\mu) &:= \langle g, \mathbf{Q}\mu \cdot \mathbf{n} \rangle_{\partial\Omega_D} + (f, \mathcal{U}\mu)_{\Omega_h} - \langle \mu, \mathbf{q}_N \rangle_{\partial\Omega_N}, \end{aligned}$$

for all η and $\mu \in M_h$.

This result allows us to shed light on the effect of local stabilization parameters τ on the approximate solution. It will also allow us to compare the RT_k , the BDM_k and the SCDG $_k$ methods; see [8] for a comparison of the hybridized version of the RT_k and the BDM_k methods. These results are gathered in the following theorem. To state it, we use the projection \mathbf{P}^{k-1} which is defined by (2.10) for $k \geq 1$ and

which we take to be identically zero when $k = 0$. We keep this convention in the remainder of the paper.

Theorem 2.4. *We have that*

- (i) *The function $(\mathbf{q}_h, \mathbf{P}^{k-1}u_h, \lambda_h)$, is independent of the values of the local stabilization parameters τ . Moreover, changes in the local stabilization parameters τ_K only affect the function $\mathcal{U}f|_{e_K^\tau}$.*
- (ii) *If $\mathbf{P}f|_K \in \mathcal{P}^{k-1}(K)$ for all simplexes $K \in \Omega_h$, then $(\mathbf{q}_h, \mathbf{P}^{k-1}u_h, \lambda_h)$ is the same for the RT_k , the BDM_k (if $k \geq 1$) and the $SCDG_k$ methods. Moreover, $u_h|_{e_K^\tau} = \hat{u}_h|_{e_K^\tau}$ for all $K \in \Omega_h$.*
- (iii) *The bilinear form $a_h(\cdot, \cdot)$ is always the same for the RT_k , the BDM_k (if $k \geq 1$) and the $SCDG_k$ methods.*

2.3. A priori error estimates. In this subsection, we obtain a priori error estimates for the error of the approximation $(\mathbf{q}_h, u_h, \lambda_h) \in \mathbf{V}_h \times W_h \times M_h$ given by the $SCDG_k$ and the numerical trace \hat{u}_h defined by (1.4a). To state them, we need to introduce new notation.

For any real-valued function ζ in $H^l(\Omega_h)$, we set

$$|\zeta|_{H^l(\Omega_h)} := \left(\sum_{K \in \Omega_h} |\zeta|_{H^l(K)}^2 \right)^{\frac{1}{2}}.$$

For a vector-valued function $\boldsymbol{\sigma} = (\sigma_1, \dots, \sigma_d) \in \mathbf{H}^l(\Omega_h)$ we set

$$|\boldsymbol{\sigma}|_{\mathbf{H}^l(\Omega_h)} := \left(\sum_{i=1}^d |\sigma_i|_{H^l(\Omega_h)}^2 \right)^{\frac{1}{2}}.$$

We can now state our results.

We begin by measuring the error in the approximation of the flux \mathbf{q} in the norm $\|\boldsymbol{\sigma}\|_{L^2(\Omega_h; \mathbf{c})} = (\mathbf{c}\boldsymbol{\sigma}, \boldsymbol{\sigma})_{\Omega_h}^{1/2}$.

Theorem 2.5. *Suppose that the exact flux \mathbf{q} belongs to $\mathbf{H}^{r+1}(\Omega_h)$ for some $r \in [0, k]$. Then*

$$\|\mathbf{q} - \mathbf{q}_h\|_{L^2(\Omega_h; \mathbf{c})} \leq \|\mathbf{q} - \mathbf{\Pi}\mathbf{q}\|_{L^2(\Omega_h; \mathbf{c})} \leq C h^{r+1} |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)},$$

for some constant C independent of h and the exact solution (\mathbf{q}, u) .

Note that the upper bound of the error is independent of the local stabilization parameters τ , in complete agreement with the characterization of the approximate solution given by Theorem 2.3. It is interesting to realize that the first estimate also holds for the RT_k and BDM_k methods when the projection $\mathbf{\Pi}$ is suitably chosen; see [11] and [4]. Such estimate is obtained by using the commutativity property and the fact that the image of their projections is in $\mathbf{H}(\text{div}, \Omega)$. Since our projection $(\mathbf{\Pi}, \mathbf{p})$ only satisfies a weak version of the commutativity property, a much more delicate analysis has to be carried out to obtain it.

In [5], it was shown that for general LDG methods with penalization parameters of order $1/h$, the order of convergence of the approximations for flux \mathbf{q} using polynomials of degree k is only k ; this order is sharp because it is actually attained for some LDG methods. It was also shown that, for DG methods with both penalization parameters of order one, the order of convergence of the approximations for the flux using polynomials of degree k is $(k + 1/2)$. Here, we obtain an order of convergence of $(k + 1)$. *No other DG method has this property.*

Next, we present several estimates for the error in the approximation of the potential u . The first is a superconvergence result. To state it, we need to introduce the adjoint equations

$$(2.13a) \quad \mathbf{c} \boldsymbol{\psi} + \nabla \varphi = 0 \quad \text{in } \Omega,$$

$$(2.13b) \quad \nabla \cdot \boldsymbol{\psi} = \theta \quad \text{in } \Omega,$$

$$(2.13c) \quad \varphi = 0 \quad \text{on } \partial\Omega_D$$

$$(2.13d) \quad \boldsymbol{\psi} \cdot \mathbf{n} = 0 \quad \text{on } \partial\Omega_N.$$

We also need to assume that the following elliptic regularity result holds

$$(2.14) \quad \|\boldsymbol{\psi}\|_{\mathbf{H}^{s+1}(\Omega_h)} + \|\nabla \varphi\|_{\mathbf{H}^{s+1}(\Omega_h)} \leq C_{er} \|\theta\|_{H^s(\Omega)}$$

for $s \in [0, k]$. Note that, since we are working with domains that can be triangulated by using straight-faced simplexes, the above result only holds if such domain is convex and $s = 0$. However, we want to write this assumption in such generality since the method will be extended to domains with smooth curved boundaries in a forthcoming paper.

Theorem 2.6. *Suppose that the exact flux \mathbf{q} belongs to $\mathbf{H}^{r+1}(\Omega_h)$ for $r \in [0, k]$. Set $\kappa := \max_{K \in \Omega_h} \frac{1}{h_K \tau_K}$. Then,*

$$\|\mathbb{P}u - u_h\|_{H^{-s}(\Omega_h)} \leq C \mathfrak{C}_\kappa^{r,s}(\mathbf{q}) h^{r+s+2},$$

where

$$\mathfrak{C}_\kappa^{r,s}(\mathbf{q}) = \begin{cases} C_{er} |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} + \kappa |\nabla \cdot \mathbf{q}|_{H^r(\Omega_h)}, & \text{for } s \in [0, k-1], k \geq 1, \\ C_{er} |\mathbf{q}|_{\mathbf{H}^1(\Omega_h)} + C_{er} |\mathbf{q}_N|_{H^1(\partial\Omega_N)} & \text{for } r = s = k = 0 \text{ and } f = 0. \end{cases}$$

Moreover, for $k = 0$ and general $f \in L^2(\Omega_h)$,

$$\|\mathbb{P}u - u_h\|_{L^2(\Omega_h)} \leq C \mathcal{C}_\kappa(\mathbf{q}) h,$$

where $\mathcal{C}_\kappa(\mathbf{q}) = (1+h)C_{er} |\mathbf{q}|_{\mathbf{H}^1(\Omega_h)} + h \kappa |\nabla \cdot \mathbf{q}|_{L^2(\Omega_h)}$.

It is interesting to note that the above superconvergence result holds for any choice of local stabilization parameters τ_K such that κ is uniformly bounded, that is, such that $1/(h_K \tau_K)$ is uniformly bounded with respect to h . This shows that τ_K cannot be too small for superconvergence to take place.

A straightforward consequence of this theorem is the following result.

Corollary 2.7. *Suppose that the exact flux (\mathbf{q}, u) belongs to $\mathbf{H}^{r+1}(\Omega_h) \times H^{r+1}(\Omega_h)$ for $r \in [0, k]$. Then*

$$\|u - u_h\|_{L^2(\Omega_h)} \leq C h^{r+1} (\mathcal{C}_\kappa(\mathbf{q}) + |u|_{H^{r+1}(\Omega_h)}),$$

where

$$\mathcal{C}_\kappa(\mathbf{q}) = \begin{cases} \min\{\mathfrak{C}_\kappa^{r-1,0}(\mathbf{q}), h \mathfrak{C}_\kappa^{r,0}(\mathbf{q})\} & \text{if } k \geq 1, \\ \mathcal{C}_\kappa(\mathbf{q}) & \text{if } k = 0. \end{cases}$$

Note that the above result shows that if $1/\tau_K$ is uniformly bounded for quasi-uniform triangulations, the convergence of u_h is still optimal, provided \mathbf{q} is smoother than required, that is, provided $\mathbf{q} \in \mathbf{H}^{r+1}(\Omega_h)$ instead of just $\mathbf{q} \in \mathbf{H}^r(\Omega_h)$. Of course, in this case, the superconvergence of u_h to $\mathbb{P}u$ is lost.

The next result is a superconvergence result for the Lagrange multiplier λ_h . To state it, we use the following norm:

$$\| \mathbf{P}_{\partial} u - \hat{u}_h \|_{L^2(\mathcal{E}_h; h)} = \left(\sum_{K \in \Omega_h} h_K \| \mathbf{P}_{\partial} u - \hat{u}_h \|_{L^2(\partial K)}^2 \right)^{1/2}.$$

Theorem 2.8. *Suppose that the exact solution (u, \mathbf{q}) of (1.1) belongs to $H^{r+1}(\Omega_h) \times \mathbf{H}^{r+1}(\Omega_h)$ for some $r \in [0, k]$. Then,*

$$\| \mathbf{P}_{\partial} u - \hat{u}_h \|_{L^2(\mathcal{E}_h; h)} \leq C \left(\mathbf{c}_0^{r,0}(\mathbf{q}) + |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} \right) h^{r+2},$$

if $k \geq 1$, or if $k = 0$ and $f = 0$.

There are no results of this type for any other DG method. However, the RT_k and the BDM_k methods have both similar results. Here we exploited the similarity of the SCDG_k methods with the RT_k and BDM_k methods to obtain these superconvergence results.

2.4. Postprocessing. We end this section by showing how to exploit the superconvergence results to postprocess u_h , \mathbf{q}_h and \hat{u}_h to get a better approximation to u defined as follows.

On the simplex K , we define the new approximation of u , u_h^* , as the function of $\mathcal{P}^{k+1}(K)$ given by

$$(2.15a) \quad u_h^* = \bar{u}_h + \tilde{u}_h,$$

where

$$(2.15b) \quad \bar{u}_h = \begin{cases} \frac{1}{d} \sum_{e \in \partial K} \hat{u}_h|_e & \text{if } k = 0, \\ \frac{1}{|K|} \int_K u_h \, dx & \text{if } k > 0, \end{cases}$$

and \tilde{u}_h is the polynomial in $\mathcal{P}_0^{k+1}(K)$ satisfying

$$(2.15c) \quad (\mathbf{a} \nabla \tilde{u}_h, \nabla w)_K = (f, w)_K - \langle w, \hat{\mathbf{q}}_h \cdot \mathbf{n} \rangle_{\partial K} \quad \forall w \in \mathcal{P}^{k+1}(K).$$

Here $\mathbf{a} = \mathbf{c}^{-1}$ and $\mathcal{P}_0^{k+1}(K)$ is the collection of functions in $\mathcal{P}^{k+1}(K)$ with mean zero. The postprocessing technique just introduced is a slight modification of a postprocessing proposed in [14, 15] and [12]; it consists in using the numerical trace $\hat{\mathbf{q}}_h$ instead of \mathbf{q}_h .

It is easy to see that this postprocessing is associated to a locally conservative method. Indeed, the scheme satisfied by u_h^* on each simplex $K \in \Omega_h$ is

$$(\mathbf{a} \nabla u_h^*, \nabla w)_K + \langle w, \hat{\mathbf{q}}_h \cdot \mathbf{n} \rangle_{\partial K} = (f, w)_K \quad \forall w \in \mathcal{P}^{k+1}(K).$$

As a consequence, if we take D_h to be the union of an arbitrary set of simplexes $K \in \Omega_h$, we get that

$$\langle 1, \hat{\mathbf{q}}_h \cdot \mathbf{n} \rangle_{\partial D_h} = (f, 1)_{D_h},$$

which is nothing by the property of local conservativity.

Note that \tilde{u}_h is well defined. Indeed, if we take $w = 1$ in equation (2.15c), the right-hand side is also equal to zero thanks to equation (1.3b). The fact that it provides a better approximation to the potential u than u_h is contained in the following result.

Theorem 2.9. *Suppose that the exact solution (u, \mathbf{q}) belongs to $H^{r+2}(\Omega_h) \times \mathbf{H}^{r+1}(\Omega_h)$ for $r \in [0, k]$. Then, if $k \geq 1$,*

$$\|u - u_h^*\|_{L^2(\Omega_h)} \leq C h^{r+2} \left(\mathbf{c}_0^{r,0}(\mathbf{q}) + |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} + |u|_{H^{r+2}(\Omega_h)} \right),$$

and if $k = 0$ and $f \equiv 0$,

$$\|u - u_h^*\|_{L^2(\Omega_h)} \leq C h^2 \left(\mathbf{c}_0(\mathbf{q}) + |u|_{H^2(\Omega_h)} \right).$$

Note that when $\mathbf{P}f|_K \in \mathcal{P}^{k-1}(K)$ for all $K \in \Omega_h$, by Theorem 2.4 we have that the function $(\hat{\mathbf{q}}_h \cdot \mathbf{n}, \mathbf{P}^{k-1}u_h, \lambda_h)$ is the same for the RT_k , BDM_k ($k \geq 1$) and SCDG_k methods. As a consequence, the postprocessed approximation u_h^* is also the same for all these methods.

Note also that in [3], a general postprocessing which is solely based on approximation results was obtained. When applied to the SCDG_k method for $k \geq 1$, it gives rise to an approximation of u which converges with the same orders as ours. However, unlike such postprocessing, our postprocessed solution u_h^* is associated to a locally conservative scheme; it is also easier to compute.

Let us end this section by noting that all the error estimates for $k \geq 1$ hold if in the equation (1.3b), we replace f by any function $\mathcal{J}_h f$ such that $\mathcal{J}_h f|_K \in \mathcal{P}^{k-1}(K)$ for all $K \in \Omega_h$ and such that

$$\|f - \mathcal{J}_h f\|_{H^{-1}(\Omega)} \leq C h^{r+1} |f|_{H^r(\Omega_h)}.$$

Moreover, by statement (ii) of Theorem 2.4, the function $(\mathbf{q}_h, \mathbf{P}^{k-1}u_h, \lambda_h)$ provided by the RT_k , the BDM_k and the SCDG_k method is the same; in particular, we have that $\mathbf{q}_h \in \mathbf{H}(\text{div}, \Omega)$. The postprocessed approximation u_h^* is also the same for those three methods.

3. PROOFS

In this section, we present detailed proofs of all our results.

3.1. Proof of Proposition 2.1: The properties of (Π, \mathbb{P}) .

3.1.1. *Two key auxiliary results about polynomials.* To prove Proposition 2.1, we begin by stating and proving two lemmas whose use is crucial in our analysis.

Lemma 3.1. *Given the face e of the simplex K and a function $z \in \mathcal{P}^k(e)$, there is a unique function $Z \in \mathcal{P}^k(K)$ such that*

$$\begin{aligned} (i) \quad \langle Z, \omega \rangle_e &= \langle z, \omega \rangle_e \quad \forall \omega \in \mathcal{P}^k(e), \\ (ii) \quad (Z, \mathbf{w})_K &= 0 \quad \forall \mathbf{w} \in \mathcal{P}^{k-1}(K). \end{aligned}$$

Moreover,

$$(iii) \quad \|Z\|_{L^2(K)} \leq C h_K^{1/2} \|z\|_{L^2(e)},$$

where h_K is the diameter of the simplex K and C depends solely on k and the shape-regularity parameters of the simplex K .

Lemma 3.2. *Given the face e of the simplex K and the function z such that for all faces e' of K different from e , $z|_{e'} \in \mathcal{P}^k(e')$, there is a unique function $\mathbf{Z} \in \mathcal{P}^k(K)$ such that*

$$\begin{aligned} (i) \quad & \langle \mathbf{Z} \cdot \mathbf{n}, \omega \rangle_{e'} = \langle z, \omega \rangle_{e'} \quad \forall \omega \in \mathcal{P}^k(e'), e' \neq e, \\ (ii) \quad & (\mathbf{Z}, \mathbf{v})_K = 0 \quad \forall \mathbf{v} \in \mathcal{P}^{k-1}(K). \end{aligned}$$

Moreover,

$$(iii) \quad \|\mathbf{Z}\|_{L^2(K)} \leq C h_K^{1/2} \|z\|_{L^2(\partial K \setminus e)},$$

where h_K is the diameter of the simplex K and C depends solely on k and the shape-regularity parameters of the simplex K .

We are only going to give a detailed proof of Lemma 3.2 since the proof of Lemma 3.1 is similar and simpler.

Proof of Lemma 3.2. Let us begin by proving that the function $\sigma \in \mathcal{P}^k(K)$ satisfying (i) and (ii) exists and is unique. Since the linear system determined by equations (i) and (ii) is square, indeed,

$$\begin{aligned} \dim(\mathcal{P}^{k-1}(K)) &= \binom{k-1+d}{d} \times d, \\ \sum_{e' \in \partial K, e' \neq e} \dim(\mathcal{P}^k(e')) &= \binom{k+d-1}{d-1} \times d, \\ \dim(\mathcal{P}^k(K)) &= \binom{k+d}{d} \times d, \end{aligned}$$

and $\binom{k-1+d}{d} + \binom{k+d-1}{d-1} = \binom{k+d}{d}$, we only need to show that if $\sigma \in \mathcal{P}^k(K)$ satisfies

$$\begin{aligned} (\mathbf{Z}, \mathbf{v})_K &= 0, \quad \forall \mathbf{v} \in \mathcal{P}^{k-1}(K), \\ \langle \mathbf{Z} \cdot \mathbf{n}, \omega \rangle_{e'} &= 0 \quad \forall \omega \in \mathcal{P}^k(e'), e' \in \partial K, e' \neq e, \end{aligned}$$

then $\mathbf{Z} = \mathbf{0}$ on K .

Let T be the affine mapping that transforms the element K to the reference simplex \widehat{K} . Moreover, let us denote by e_i , $i = 1, \dots, d+1$, the faces of K where $e := e_{d+1}$. Assume that the mapping T is such that $\widehat{e}_i := T(e_i)$ is the face of \widehat{K} lying on the plane $\widehat{x}_i = 0$. and set $\widehat{\mathbf{Z}}(\widehat{x}) := \mathbf{Z}(T^{-1}(\widehat{x}))$. Then the above equations become

$$\begin{aligned} (\widehat{\mathbf{Z}}, \widehat{\mathbf{v}})_{\widehat{K}} &= 0 \quad \forall \widehat{\mathbf{v}} \in \mathcal{P}^{k-1}(\widehat{K}), \\ \langle \widehat{\mathbf{Z}} \cdot \mathbf{n}_i, \widehat{\omega} \rangle_{\widehat{e}_i} &= 0 \quad \forall \widehat{\omega} \in \mathcal{P}^k(\widehat{e}_i), i = 1, \dots, d, \end{aligned}$$

since spaces of polynomials of a given total degree are invariant under affine transformations. Now, let $\{\widehat{\mathbf{n}}_j\}_{j=1}^d$ be the basis of \mathbb{R}^d dual to $\{\mathbf{n}_i\}_{i=1}^d$, that is, $\mathbf{n}_i \cdot \widehat{\mathbf{n}}_j = \delta_{ij}$. Then we can write $\widehat{\mathbf{Z}} = \sum_{j=1}^d \widehat{p}_j \widehat{\mathbf{n}}_j$, where $\widehat{p}_j \in \mathcal{P}^k(\widehat{K})$, and obtain that

$$\begin{aligned} \sum_{j=1}^d (\widehat{p}_j, \widehat{\mathbf{n}}_j \cdot \widehat{\mathbf{v}})_{\widehat{K}} &= 0 \quad \forall \widehat{\mathbf{v}} \in \mathcal{P}^{k-1}(\widehat{K}), \\ \langle \widehat{p}_i, \widehat{\omega} \rangle_{\widehat{e}_i} &= 0 \quad \forall \widehat{\omega} \in \mathcal{P}^k(\widehat{e}_i), i = 1, \dots, d. \end{aligned}$$

The last equation implies that, for any $i = 1, \dots, d$, $\widehat{p}_i|_{\widehat{e}_i} = 0$ and hence that $\widehat{p}_i = \widehat{x}_i \widehat{\mathbf{p}}_i$ for some polynomial $\widehat{\mathbf{p}}_i$ in $\mathcal{P}^{k-1}(\widehat{K})$. Taking $\widehat{\mathbf{v}} = \widehat{\mathbf{p}}_i \mathbf{n}_i$, we get

$$(\widehat{p}_i, \widehat{\mathbf{p}}_i)_{\widehat{K}} = (\widehat{x}_i, \widehat{\mathbf{p}}_i^2)_{\widehat{K}} = 0,$$

and, since $\widehat{x}_i > 0$ on \widehat{K} , we conclude that $\widehat{\mathbf{p}}_i = 0$. This implies that $\mathbf{Z} = \mathbf{0}$ on K . This proves the existence and uniqueness of \mathbf{Z} satisfying the conditions (i) and (ii).

The estimate (iii) follows now from a simple scaling argument. This completes the proof. \square

3.1.2. Proof of the orthogonality properties. It is not difficult to see that the fact that $(\mathbf{\Pi}, \mathbb{P})$ is well defined is a direct corollary of Lemmas 3.1 and 3.2.

Now, let us prove the orthogonality properties. The property (i) follows from the property (2.8a) defining \mathbb{P} and the orthogonality property (ii) follows from the property (2.7a) defining $\mathbf{\Pi}$. The orthogonality property (iii) follows from the properties (2.8b) and (2.7b) defining \mathbb{P} and $\mathbf{\Pi}$, and from the definition of the projection \mathbb{P}_∂ , (1.5). In fact, it follows from the fact that on each face e of any simplex K , we have that either $\mathbb{P}\zeta = \mathbb{P}_\partial\zeta$ or $\mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} = \mathbb{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n}$.

3.1.3. Proof of the weak commutativity property. The weak commutativity property (iv) is a direct consequence of the three orthogonality properties we just proved. Indeed, we have that

$$\begin{aligned} -(\nabla\zeta, \mathbf{\Pi}\boldsymbol{\sigma})_K &= (\zeta, \nabla \cdot \mathbf{\Pi}\boldsymbol{\sigma})_K - \langle \zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K} \\ &= (\mathbb{P}\zeta, \nabla \cdot \mathbf{\Pi}\boldsymbol{\sigma})_K - \langle \zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K} && \text{by (i),} \\ &= -(\nabla\mathbb{P}\zeta, \mathbf{\Pi}\boldsymbol{\sigma})_K + \langle \mathbb{P}\zeta - \zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K} \\ &= -(\nabla\mathbb{P}\zeta, \boldsymbol{\sigma})_K + \langle \mathbb{P}\zeta - \zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K} && \text{by (ii),} \\ &= -(\nabla\mathbb{P}\zeta, \boldsymbol{\sigma})_K + \langle \mathbb{P}\zeta - \mathbb{P}_\partial\zeta, \boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K} && \text{by (iii),} \\ &= (\mathbb{P}\zeta, \nabla \cdot \boldsymbol{\sigma})_K - \langle \mathbb{P}_\partial\zeta, \boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K}, \\ &= (\mathbb{P}\zeta, \nabla \cdot \boldsymbol{\sigma})_K - \langle \mathbb{P}_\partial\zeta, \mathbb{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial K}, \end{aligned}$$

by the definition of the projection \mathbb{P}_∂ , (1.5). This completes the proof of (iv).

3.1.4. Proof of the estimates (v) and (vi). Note that, by the definition of the projections $\mathbf{\Pi}$, (2.7), and $\mathbf{\Pi}^{\text{RT}}$, (2.9), we have that

$$\begin{aligned} (\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma} - \mathbf{\Pi}\boldsymbol{\sigma}, \mathbf{v})_K &= 0, \\ \langle (\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma} - \mathbf{\Pi}\boldsymbol{\sigma}) \cdot \mathbf{n}, \omega \rangle_e &= \langle \mathbb{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n}, \omega \rangle_{e \cap e_K^\tau}. \end{aligned}$$

for all $\mathbf{v} \in \mathcal{P}^{k-1}(K)$, if $k > 1$, and for all $\omega \in \mathcal{P}^k(e)$ and all faces e of K . By a well known scaling argument, we immediately obtain that

$$\|\mathbf{\Pi}^{\text{RT}}\boldsymbol{\sigma} - \mathbf{\Pi}\boldsymbol{\sigma}\|_{L^2(K)} \leq C h_K^{1/2} \|\mathbb{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n}\|_{L^2(e_K^\tau)}.$$

It remains to estimate the above right-hand side. To do that, we note that, for any ω in $\mathcal{P}^k(K)$, we have that

$$\langle \omega, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbb{P}_\partial\boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{e_K^\tau} = \langle \omega, (\mathbf{\Pi}\boldsymbol{\sigma} - \boldsymbol{\sigma}) \cdot \mathbf{n} \rangle_{\partial K} = (\omega, \nabla \cdot (\mathbf{\Pi}\boldsymbol{\sigma} - \boldsymbol{\sigma}))_K,$$

by the definition of the projections $\mathbf{\Pi}$, (2.7). Taking $\omega = Z$, where Z is given by Lemma 3.1 with $z := \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{P}_{\partial}\boldsymbol{\sigma} \cdot \mathbf{n}$, we get that, for any p in $\mathcal{P}^{k-1}(K)$,

$$\begin{aligned} \|z\|_{L^2(e_K^{\tau})}^2 &= \|\mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{P}_{\partial}\boldsymbol{\sigma} \cdot \mathbf{n}\|_{L^2(e_K^{\tau})}^2 \\ &= -(\mathbf{P}\nabla \cdot \boldsymbol{\sigma} - p, Z)_K \\ &\leq C h_K^{1/2} \|\mathbf{P}\nabla \cdot \boldsymbol{\sigma} - p\|_{L^2(K)} \|z\|_{L^2(e_K^{\tau})}, \end{aligned}$$

and, after a direct application of the Bramble-Hilbert lemma, we get

$$\|\mathbf{P}_{\partial}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n}\|_{L^2(e_K^{\tau})} \leq C h_K^{r+1/2} |\mathbf{P}\nabla \cdot \boldsymbol{\sigma}|_{H^r(K)},$$

where $r \in [0, k]$. This completes the proof of the estimates (v) and (vi).

3.1.5. *Proof of the estimate (vii).* Note that, by the definition of the projections \mathbb{P} , (2.8), and \mathbf{P} , (2.10), we have that

$$\begin{aligned} (\mathbb{P}\zeta - \mathbf{P}\zeta, \mathbf{w})_K &= 0, \\ \langle \mathbb{P}\zeta - \mathbf{P}\zeta, \boldsymbol{\omega} \rangle_{e_K^{\tau}} &= \langle \mathbf{P}_{\partial}\zeta - \mathbf{P}\zeta, \boldsymbol{\omega} \rangle_{e_K^{\tau}}, \end{aligned}$$

for all $\mathbf{w} \in \mathcal{P}^{k-1}(K)$, if $k \geq 1$, and for all $\boldsymbol{\omega} \in \mathcal{P}^k(e_K^{\tau})$. This implies that Lemma 3.1 holds with $z := \mathbf{P}_{\partial}\zeta - \mathbf{P}\zeta$ and $Z = \mathbb{P}\zeta - \mathbf{P}\zeta$. As a consequence,

$$\|\mathbb{P}\zeta - \mathbf{P}\zeta\|_{L^2(K)} \leq C h_K^{1/2} \|\mathbf{P}_{\partial}\zeta - \mathbf{P}\zeta\|_{L^2(e_K^{\tau})}.$$

It remains to estimate the above right-hand side.

To do that, we note that, for any \mathbf{v} in $\mathcal{P}^k(K) \oplus \boldsymbol{x} \mathcal{P}^k(K)$, we have that

$$(3.16) \quad \langle \mathbf{P}\zeta - \mathbf{P}_{\partial}\zeta, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} = \langle \mathbf{P}\zeta - \zeta, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} = (\nabla \mathbf{P}\zeta - \nabla \zeta, \mathbf{v})_K,$$

by the definition of the projection \mathbf{P} , (2.10) and that of the projection \mathbf{P}_{∂} , (1.5). A well known scaling argument states that given any function z such that its restriction to each face e of k belongs to $\mathcal{P}^k(e)$, there is a function \mathbf{Z} in $\mathcal{P}^k(K) \oplus \boldsymbol{x} \mathcal{P}^k(K)$ such that

$$\begin{aligned} (i) \quad \langle \mathbf{Z} \cdot \mathbf{n}, \boldsymbol{\omega} \rangle_e &= \langle z, \boldsymbol{\omega} \rangle_e \quad \forall \boldsymbol{\omega} \in \mathcal{P}^k(e), \\ (ii) \quad (\mathbf{Z}, \mathbf{v})_K &= 0 \quad \forall \mathbf{v} \in \mathcal{P}^{k-1}(K). \\ (iii) \quad \|\mathbf{Z}\|_{L^2(K)} &\leq C h_K^{1/2} \|z\|_{L^2(\partial K)}, \end{aligned}$$

where h_K is the diameter of the simplex K and C depends solely on the shape-regularity constants of the simplex K . Taking $\mathbf{v} := \mathbf{Z}$ with $z = \mathbf{P}\zeta - \mathbf{P}_{\partial}\zeta$ in equation (3.16), we obtain that

$$\begin{aligned} \|z\|_{L^2(\partial K)}^2 &= (\nabla \mathbf{P}\zeta - \nabla \zeta, \mathbf{Z})_K \\ &= (\mathbf{p} - \nabla \zeta, \mathbf{Z})_K \\ &\leq C h_K^{1/2} \|\nabla \zeta - \mathbf{p}\|_{L^2(K)} \|z\|_{L^2(\partial K)}, \end{aligned}$$

for any $\mathbf{p} \in \mathcal{P}^{k-1}(K)$. Thus, after a direct application of the Bramble-Hilbert lemma, we get

$$\|\mathbf{P}\zeta - \mathbf{P}_{\partial}\zeta\|_{L^2(\partial K)} \leq C h_K^{r+1/2} |\nabla \zeta|_{H^r(K)},$$

where $r \in [0, k]$. This completes the proof of estimate (vii).

This completes the proof of Proposition 2.1.

3.2. Characterization of the approximate solution. To prove the results about the characterization of the approximate solution of the SCDG method, we begin by proving two auxiliary results concerning key properties of the local solvers.

3.2.1. Two auxiliary results about the local solvers. To state the first auxiliary result, we need to introduce the following decomposition of our local spaces:

$$\begin{aligned}\mathcal{P}^k(K) &= \mathcal{V}(K) \oplus \mathcal{V}^\perp(K), \\ \mathcal{P}^k(K) &= \mathcal{W}(K) \oplus \mathcal{W}^\perp(K),\end{aligned}$$

where

$$\begin{aligned}\mathcal{V}(K) &:= \{\mathbf{v} \in \mathcal{P}^k(K) : \nabla \cdot \mathbf{v} = 0\}, \\ \mathcal{V}^\perp(K) &:= \{\mathbf{v} \in \mathcal{P}^k(K) : (\mathbf{c} \mathbf{v}, \boldsymbol{\sigma})_K = 0 \ \forall \boldsymbol{\sigma} \in \mathcal{V}(K)\},\end{aligned}$$

and

$$\begin{aligned}\mathcal{W}(K) &:= \mathcal{P}^{k-1}(K), \\ \mathcal{W}^\perp(K) &:= \{\omega \in \mathcal{P}^k(K) : (\omega, \zeta)_K = 0 \ \forall \zeta \in \mathcal{W}(K)\},\end{aligned}$$

Lemma 3.3. *Let K be any simplex of the triangulation Ω_h . Then the local mapping $(\mathbf{Qm}, \mathcal{Um})$ given by equations (2.11) can be obtained as follows:*

(i) *Set*

$$\mathcal{Um}|_{e_K^\tau} = \mathbf{P}_{\partial\mathbf{m}}|_{e_K^\tau}.$$

(ii) *Compute $\mathbf{Qm} \in \mathcal{V}(K)$ by solving*

$$(\mathbf{c} \mathbf{Qm}, \mathbf{v})_K = -\langle \mathbf{m}, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} \quad \forall \mathbf{v} \in \mathcal{V}(K).$$

(iii) *Compute $\mathbf{P}^{k-1} \mathcal{Um}$ by solving*

$$(\mathbf{P}^{k-1} \mathcal{Um}, \nabla \cdot \mathbf{v})_K = \langle \mathbf{m}, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} \quad \forall \mathbf{v} \in \mathcal{V}^\perp(K).$$

Similarly, the local mapping $(\mathbf{Q}f, \mathcal{U}f)$ given by equations (2.12) can be obtained as follows:

(α) *Compute $\mathcal{U}f|_{e_K^\tau}$ by solving*

$$\langle \omega, \tau \mathcal{U}f \rangle_{e_K^\tau} = (f, \omega)_K \quad \forall \omega \in \mathcal{W}^\perp(K).$$

(β) *Compute $\mathbf{Q}f \in \mathcal{V}^\perp(K)$ by solving*

$$(\omega, \nabla \cdot \mathbf{Q}f)_K = -\langle \omega, \tau \mathcal{U}f \rangle_{e_K^\tau} + (f, \omega)_K \quad \forall \omega \in \mathcal{W}(K).$$

(γ) *Compute $\mathbf{P}^{k-1} \mathcal{U}f$ by solving*

$$(\mathbf{P}^{k-1} \mathcal{U}f, \nabla \cdot \mathbf{v})_K = (\mathbf{c} \mathbf{Q}f, \mathbf{v})_K \quad \forall \mathbf{v} \in \mathcal{V}^\perp(K).$$

Proof. Let us begin by proving the properties of the first local mapping. Thus, integrating by parts in the equation (2.11b), we obtain

$$(\omega, \nabla \cdot \mathbf{Qm})_K = \langle \omega, (\widehat{\mathbf{Qm}} - \mathbf{Qm}) \cdot \mathbf{n} \rangle_{\partial K} = \langle \omega, \tau(\mathcal{Um} - \mathbf{m}) \rangle_{e_K^\tau}$$

for all $\omega \in \mathcal{P}^k(K)$, by the definition of the numerical trace $\widehat{\mathbf{Qm}}$, (1.4b) and (1.6). Taking $\omega \in \mathcal{W}^\perp(K)$, we see that

$$\langle \omega, \tau(\mathcal{Um} - \mathbf{m}) \rangle_{e_K^\tau} = 0 \quad \forall \omega \in \mathcal{W}^\perp(K).$$

Using the fact that $\mathcal{W}^\perp(K)|_{e_K^\tau} = \mathcal{P}^k(e_K^\tau)$ which follows by a simple application of Lemma (3.1), we have that (i) holds. As a consequence, we see that

$$(\omega, \nabla \cdot \mathbf{Qm})_K = 0 \quad \forall \omega \in \mathcal{P}^k(K),$$

and hence that $\mathbf{Qm} \in \mathbf{V}(K)$. The property (ii) can be now obtained by restricting the test functions \mathbf{v} to the space $\mathbf{V}(K)$ in the equation (2.11a). Now that we know \mathbf{Qm} , we obtain the formulation (iii) for $\mathcal{P}^{k-1}\mathcal{U}\mathbf{m}$ by restricting the test functions \mathbf{v} to the space $\mathbf{V}^\perp(K)$ in the equation (2.11a). It remains to show that $\mathcal{P}^{k-1}\mathcal{U}\mathbf{m}$ is uniquely defined by those equations. But this follows from the fact that the system of equations is square and that $\nabla \cdot \mathbf{V}^\perp(K) = \mathcal{P}^{k-1}(K) = \mathcal{W}(K)$, which in turn follows from the fact that $\nabla \cdot \mathbf{V}(K) = \{0\}$ and $\nabla \cdot \mathcal{P}^k(K) = \mathcal{P}^{k-1}(K)$. This completes the proof of the properties of the first local lifting.

The proof the properties (α) and (γ) of the second local mapping is similar to the proof of the properties (i) and (iii) of the first local mapping, respectively. Let us prove property (β) . If we take $\mathbf{v} \in \mathbf{V}$ in the equation (2.12a), we see that $\mathbf{Q}f \in \mathbf{V}^\perp(K)$. Since the equation in (β) is obtained from (2.12b) by restricting the tests functions ω to $\mathcal{W}(K)$, we only have to prove that $\mathbf{Q}f$ given by (β) is well defined. But this follows from the fact that the system is a square system and $\nabla \cdot \mathbf{V}^\perp(K) = \mathcal{P}^{k-1}(K) = \mathcal{W}(K)$. This completes the proof. \square

The second auxiliary result concerns the jumps of the local solvers.

Lemma 3.4. *For each simplex $K \in \Omega_h$, we have that, on ∂K ,*

$$\begin{aligned} (\widehat{\mathbf{Qm}} - \mathbf{Qm}) \cdot \mathbf{n} &= \tau (\mathcal{U}\mathbf{m} - \mathcal{P}_\partial \mathbf{m}) = 0, \\ (\widehat{\mathbf{Q}f} - \mathbf{Q}f) \cdot \mathbf{n} &= \tau \mathcal{U}f = \mathcal{P}_\partial \mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi} \mathbf{q} \cdot \mathbf{n}. \end{aligned}$$

Proof. Let us begin by proving the second identity since its proof is more involved. Taking $\omega = Z$ in the identity (β) of Lemma 3.3, where Z is given by Lemma 3.1 with $e = e_K^\tau$, we obtain that

$$\begin{aligned} \langle z, (\widehat{\mathbf{Q}f} - \mathbf{Q}f) \cdot \mathbf{n} \rangle_{e_K^\tau} &= \langle z, \tau \mathcal{U}f \rangle_{e_K^\tau} \\ &= (Z, \nabla \cdot (\mathbf{q} - \mathbf{\Pi} \mathbf{q}))_K \\ &= -(\nabla Z, \mathbf{q} - \mathbf{\Pi} \mathbf{q})_K + \langle Z, (\mathbf{q} - \mathbf{\Pi} \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial K} \\ &= \langle z, \mathcal{P}_\partial \mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi} \mathbf{q} \cdot \mathbf{n} \rangle_{e_K^\tau}, \end{aligned}$$

by the properties of the projection $\mathbf{\Pi}$, (2.7). As a consequence, we immediately obtain that, on e_K^τ ,

$$(\widehat{\mathbf{Q}f} - \mathbf{Q}f) \cdot \mathbf{n} = \tau_K \mathcal{U}f = \mathcal{P}_\partial \mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi} \mathbf{q} \cdot \mathbf{n}.$$

A similar argument gives that, on e_K^τ ,

$$(\widehat{\mathbf{Qm}} - \mathbf{Qm}) \cdot \mathbf{n} = \tau_K (\mathcal{U}\mathbf{m} - \mathcal{P}_\partial \mathbf{m}) = 0.$$

This completes the proof of Lemma 3.4. \square

3.2.2. Proof of Theorem 2.3: Characterization of the approximate solution. The following result is a particular case of a general result for hybridized Galerkin methods proven in [9].

Theorem 3.5. *The approximate solution $(\mathbf{q}_h, u_h, \lambda_h) \in \mathbf{V}_h \times W_h \times M_h$ given by the SCDG $_k$ method is well defined. Moreover, we have that*

$$(\mathbf{q}_h, u_h) = (\mathbf{Q}\lambda_h, \mathcal{U}\lambda_h) + (\mathbf{Q}g, \mathcal{U}g) + (\mathbf{Q}f, \mathcal{U}f),$$

where λ_h can be characterized as the function in M_h satisfying

$$a_h(\lambda_h, \mu) = b_h(\mu) \quad \forall \mu \in M_h,$$

where

$$\begin{aligned} a_h(\eta, \mu) &= (c\mathbf{Q}\eta, \mathbf{Q}\mu)_{\Omega_h} - \langle \mu - \mathcal{U}\mu, (\widehat{\mathbf{Q}}\eta - \mathbf{Q}\eta) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ b_h(\mu) &= \langle g, \widehat{\mathbf{Q}}\mu \cdot \mathbf{n} \rangle_{\partial\Omega_D} + (f, \mathcal{U}\mu)_{\Omega_h} - \langle \mu, \mathbf{q}_N \rangle_{\partial\Omega_N} \\ &\quad - \langle \mu - \mathcal{U}\mu, (\widehat{\mathbf{Q}}f - \mathbf{Q}f) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad + \langle -\mathcal{U}f, (\widehat{\mathbf{Q}}\mu - \mathbf{Q}\mu) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad + \langle \mu - \mathcal{U}\mu, (\widehat{\mathbf{Q}}g - \mathbf{Q}g) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad - \langle g - \mathcal{U}g, (\widehat{\mathbf{Q}}\mu - \mathbf{Q}\mu) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \end{aligned}$$

for all η and $\mu \in M_h$.

Theorem 2.3 follows from this result if we show that on \mathcal{E}_h ,

$$(\widehat{\mathbf{Q}}\mu - \mathbf{Q}\mu) \cdot \mathbf{n} = 0 \quad \text{and} \quad (\mu - \mathcal{U}\mu)(\widehat{\mathbf{Q}}f - \mathbf{Q}f) = 0,$$

for all $\mu \in M_h$. Since this is a straightforward consequence of Lemma 3.4, this completes the proof of Theorem 2.3.

3.2.3. Proof of Proposition 2.2: Characterization of the jumps. By the definition of the numerical traces (1.4) and (1.6), we have that

$$\begin{aligned} (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} &= \begin{cases} \tau_K (u_h - \widehat{u}_h) & \text{on } e_K^{\tau}, \\ 0 & \text{otherwise.} \end{cases} \\ &= \begin{cases} \tau_K (u_h - \widehat{u}_h) & \text{on } e_K^{\tau}, \\ \mathbf{P}_{\partial}\mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} & \text{otherwise.} \end{cases} \end{aligned}$$

by the definition of the projection $\mathbf{\Pi}$, (2.7b), and that of the projection \mathbf{P}_{∂} , (1.5). So, we only have to prove that

$$(\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} = \mathbf{P}_{\partial}\mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} \quad \text{on } e_K^{\tau}.$$

But, by Theorem 2.3, we have that

$$\begin{aligned} (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} &= (\widehat{\mathbf{Q}}\lambda_h - \mathbf{Q}\lambda_h) \cdot \mathbf{n} + (\widehat{\mathbf{Q}}g - \mathbf{Q}g) \cdot \mathbf{n} + (\widehat{\mathbf{Q}}f - \mathbf{Q}f) \cdot \mathbf{n} \\ &= \mathbf{P}_{\partial}\mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} \end{aligned}$$

on the face e_K^{τ} . This completes the proof of Proposition 2.2.

3.2.4. Proof of Theorem 2.4. The statement (i) of Theorem 2.4 follows directly from Theorem 2.3 and from Lemma 3.3.

To prove the remaining statements, we are going to use the fact that the RT $_k$, BDM $_k$ and SCDG $_k$ methods have exactly the same structure and satisfy the characterization theorem 2.3; see [9]. The only difference between these methods is the choice of local spaces, see Table 1, and the choice of the local stabilization parameters τ , see Table 2. Thus, to prove statement (ii) we only have to show that the functions $(\mathbf{Q}\mathbf{m}, \mathbf{P}^{k-1}\mathcal{U}\mathbf{m})$ and $(\mathbf{Q}f, \mathbf{P}^{k-1}\mathcal{U}f)$ are the same for all these methods

whenever $f|_K \in \mathcal{P}^{k-1}(K)$ for all $K \in \Omega_h$. Similarly, to prove statement (iii), we only have to show that $\mathcal{Q}\mathbf{m}$ is the same for all these methods.

To do that, we begin by noting that we have that, by Lemma 3.3, the function $(\mathcal{Q}\mathbf{m}, \mathcal{P}^{k-1}\mathcal{U}\mathbf{m}) \in \mathcal{V}(K) \times \mathcal{W}(K)$ is determined by

$$\begin{aligned} (\mathbf{c}\mathcal{Q}\mathbf{m}, \mathbf{v})_K &= -\langle \mathbf{m}, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} & \forall \mathbf{v} \in \mathcal{V}(K), \\ (\mathcal{P}^{k-1}\mathcal{U}\mathbf{m}, \nabla \cdot \mathbf{v})_K &= (\mathbf{c}\mathcal{Q}\mathbf{m}, \mathbf{v})_K + \langle \mathbf{m}, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K} & \forall \mathbf{v} \in \mathcal{V}^\perp(K). \end{aligned}$$

and $(\mathcal{Q}f, \mathcal{P}^{k-1}\mathcal{U}f) \in \mathcal{V}^\perp(K) \times \mathcal{W}(K)$ is determined by the equations

$$\begin{aligned} (\omega, \nabla \cdot \mathcal{Q}f)_K &= -\langle \omega, \tau \mathcal{U}f \rangle_{e_K^*} + (f, \omega)_K & \forall \omega \in \mathcal{W}(K), \\ (\mathcal{P}^{k-1}\mathcal{U}f, \nabla \cdot \mathbf{v})_K &= (\mathbf{c}\mathcal{Q}\mathbf{m}, \mathbf{v})_K & \forall \mathbf{v} \in \mathcal{V}^\perp(K), \end{aligned}$$

where $\mathcal{U}f|_{e_K^*} = 0$, by (α) of Lemma 3.3, if $f|_K \in \mathcal{P}^{k-1}(K)$. Since the four above equations also hold (the third whenever $f|_K \in \mathcal{P}^{k-1}(K)$) for the BDM $_k$ method, we conclude that the statements (ii) and (iii) hold if we exclude the RT $_k$ method.

To show that these statements also hold if we include it, we note that the above equations hold for the RT $_k$ method if we modify the definition of the spaces $\mathcal{V}(K)$ and $\mathcal{V}^\perp(K)$ by

$$\begin{aligned} \mathcal{V}_{\text{RT}}(K) &:= \{\mathbf{v} \in \mathcal{P}^k(K) \oplus \mathbf{x}\mathcal{P}^k(K) : \nabla \cdot \mathbf{v} = 0\}, \\ \mathcal{V}_{\text{RT}}^\perp(K) &:= \{\mathbf{v} \in \mathcal{P}^k(K) \oplus \mathbf{x}\mathcal{P}^k(K) : (\mathbf{c}\mathbf{v}, \boldsymbol{\sigma})_K = 0 \ \forall \boldsymbol{\sigma} \in \mathcal{V}_{\text{RT}}(K)\}, \end{aligned}$$

and if we replace the third equation by

$$\nabla \cdot \mathcal{Q}f = \mathcal{P}f.$$

Thus, the result follows from the fact that

$$\mathcal{V}_{\text{RT}}(K) = \{\mathbf{v} \in \mathcal{P}^k(K) : \nabla \cdot \mathbf{v} = 0\} = \mathcal{V}(K),$$

and from the fact that, if $\mathcal{Q}f \in \mathcal{V}_{\text{RT}}^\perp(K)$ and $\nabla \cdot \mathcal{Q}f = \mathcal{P}f \in \mathcal{P}^{k-1}(K)$, then $\mathcal{Q}f$ belongs to the space $\{\mathbf{v} \in \mathcal{P}^k(K) : (\mathbf{c}\mathbf{v}, \boldsymbol{\sigma})_K = 0 \ \forall \boldsymbol{\sigma} \in \mathcal{V}(K)\} = \mathcal{V}^\perp(K)$. This completes the proof of Theorem 2.4.

3.3. Proof of the error estimates. The proof of the error estimates is based on the so-called error equations and the properties of the projection $(\mathbf{\Pi}, \mathbb{P})$ gathered in Proposition 2.1. The error equations are

$$(3.17a) \quad (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{v})_{\Omega_h} - (u - u_h, \nabla \cdot \mathbf{v})_{\Omega_h} + \langle u - \hat{u}_h, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial\Omega_h} = 0,$$

$$(3.17b) \quad (\omega, \nabla \cdot (\mathbf{q} - \mathbf{q}_h))_{\Omega_h} - \langle \omega, (\hat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h} = 0,$$

$$(3.17c) \quad u - \hat{u}_h = g - \mathcal{P}_\partial g \quad \text{on } \partial\Omega_D,$$

$$(3.17d) \quad (\mathbf{q} - \hat{\mathbf{q}}_h) \cdot \mathbf{n} = \mathbf{q}_N - \mathcal{P}_\partial \mathbf{q}_N \quad \text{on } \partial\Omega_N,$$

for all $(\mathbf{v}, \omega) \in \mathcal{V}_h \times W_h$.

A direct consequence of the weak commutativity identity (iv) of Proposition 2.1 we find convenient to use in our analysis is contained in the following result.

Corollary 3.6. *For all $(\boldsymbol{\sigma}, \zeta) \in \mathbf{H}^1(\Omega_h) \times H^1(\Omega_h)$, we have*

$$\begin{aligned} (\alpha) \quad & (\zeta, \nabla \cdot \mathbf{\Pi}\boldsymbol{\sigma})_{\Omega_h} = (\mathbb{P}\zeta, \nabla \cdot \boldsymbol{\sigma})_{\Omega_h} + \langle \mathcal{P}_\partial \zeta, \mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} - \mathcal{P}_\partial \boldsymbol{\sigma} \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ (\beta) \quad & \|\mathbf{\Pi}\boldsymbol{\sigma} - \boldsymbol{\sigma}\|_{L^2(\Omega_h)} \leq C h^{r+1} |\boldsymbol{\sigma}|_{\mathbf{H}^{r+1}(\Omega_h)}, \\ (\gamma) \quad & \|\mathbb{P}\zeta - \zeta\|_{L^2(\Omega_h)} \leq C h^{s+1} |\zeta|_{H^{s+1}(\Omega_h)}, \end{aligned}$$

where $r, s \in [0, k]$ and C depends only on k and the shape-regularity parameters of the simplexes $K \in \Omega_h$.

3.3.1. *Proof of Theorem 2.5: The error in the flux.* Theorem 2.5 follows immediately from the following auxiliary result.

Lemma 3.7. *We have $(\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi q} - \mathbf{q}_h)_{\Omega_h} = 0$.*

Indeed, this implies that

$$(\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{q} - \mathbf{q}_h)_{\Omega_h} = (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{q} - \mathbf{\Pi q})_{\Omega_h},$$

and hence, that

$$\|\mathbf{q} - \mathbf{q}_h\|_{\mathbf{L}^2(\Omega_h, \mathbf{c})} \leq \|\mathbf{\Pi q} - \mathbf{q}\|_{\mathbf{L}^2(\Omega_h, \mathbf{c})} \leq Ch^{s+1} \|\mathbf{q}\|_{\mathbf{H}^{s+1}(\Omega_h)},$$

for some $s \in [0, k]$, by the estimate (vi) of Proposition 2.1. This proves Theorem 2.5.

Let us prove Lemma 3.7.

Proof. By the error equation (3.17a) with $\mathbf{v} := \mathbf{\Pi q} - \mathbf{q}_h$, we have

$$\begin{aligned} (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi q} - \mathbf{q}_h)_{\Omega_h} &= (u - u_h, \nabla \cdot (\mathbf{\Pi q} - \mathbf{q}_h))_{\Omega_h} \\ &\quad - \langle u - \hat{u}_h, (\mathbf{\Pi q} - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}. \end{aligned}$$

By the identity (α) of Corollary 3.6 with $(\boldsymbol{\sigma}, \zeta) := (\mathbf{q} - \mathbf{q}_h, u - u_h)$, we get that

$$\begin{aligned} (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi q} - \mathbf{q}_h)_{\Omega_h} &= (\mathbb{P}u - u_h, \nabla \cdot (\mathbf{q} - \mathbf{q}_h))_{\Omega_h} \\ &\quad + \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi q} - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad - \langle u - \hat{u}_h, (\mathbf{\Pi q} - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \end{aligned}$$

and by the error equation (3.17b) with $\omega := \mathbb{P}u - u_h$,

$$\begin{aligned} (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi q} - \mathbf{q}_h)_{\Omega_h} &= \langle \mathbb{P}u - u_h, (\hat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad + \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi q} - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad - \langle u - \hat{u}_h, (\mathbf{\Pi q} - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}. \end{aligned}$$

It we denote by T the right-hand side of the above equations, it is not difficult to see that, after a few simple algebraic manipulations, we have that $T = \sum_{i=1}^5 T_i$, where

$$\begin{aligned} T_1 &:= \langle \hat{u}_h - u_h, (\hat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_2 &:= \langle \hat{u}_h - u_h, (\mathbf{\Pi q} - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_3 &:= \langle \hat{u}_h - u, (\mathbf{q} - \hat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_4 &:= \langle \mathbb{P}_{\partial}u - u, (\mathbf{\Pi q} - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_5 &:= \langle \mathbb{P}u - u, (\hat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_h}. \end{aligned}$$

We are going to show that $T = 0$.

We begin by noting that,

$$\begin{aligned} T_1 &= \sum_{K \in \Omega_h} \langle \hat{u}_h - u_h, (\hat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial K} \\ &= \sum_{K \in \Omega_h} \langle \hat{u}_h - u_h, \mathbb{P}_{\partial} \mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi q} \cdot \mathbf{n} \rangle_{\partial K}, \end{aligned}$$

by Proposition 2.2. By the definition of the projection \mathbf{P}_∂ , (1.5),

$$\begin{aligned} T_1 &= \sum_{K \in \Omega_h} \langle \widehat{u}_h - u_h, (\mathbf{q} - \mathbf{\Pi}\mathbf{q}) \cdot \mathbf{n} \rangle_{\partial K} \\ &= -T_2, \end{aligned}$$

by the definition of the projection $\mathbf{\Pi}$, (2.7b). Thus, $T_1 + T_2 = 0$.

Next, let us show that $T_3 + T_4 = 0$. By the fact that the numerical trace \widehat{u}_h and the normal component of the numerical trace $\widehat{\mathbf{q}}_h$ are single-valued on the interior faces, by definition of \widehat{u}_h , (1.4a), and the equation (1.3c) satisfied by $\widehat{\mathbf{q}}_h$, we have that

$$\begin{aligned} T_3 &= \langle \widehat{u}_h - u, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial\Omega}, \\ &= \langle \mathbf{P}_\partial g - g, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial\Omega_D} + \langle \lambda_h - u, \mathbf{q}_N - \mathbf{P}_\partial \mathbf{q}_N \rangle_{\partial\Omega_N}, \end{aligned}$$

by the definition of the numerical traces at the boundary. By using the definition of the projection \mathbf{P}_∂ , (1.5), we get

$$\begin{aligned} T_3 &= \langle \mathbf{P}_\partial u - u, \mathbf{q} \cdot \mathbf{n} - \mathbf{P}_\partial \mathbf{q} \cdot \mathbf{n} \rangle_{\partial\Omega_D} + \langle \mathbf{P}_\partial u - u, \mathbf{q} \cdot \mathbf{n} - \mathbf{P}_\partial \mathbf{q} \cdot \mathbf{n} \rangle_{\partial\Omega_N} \\ &= \langle \mathbf{P}_\partial u - u, \mathbf{q} \cdot \mathbf{n} - \mathbf{P}_\partial \mathbf{q} \cdot \mathbf{n} \rangle_{\partial\Omega} \\ &= \langle \mathbf{P}_\partial u - u, \mathbf{q} \cdot \mathbf{n} - \mathbf{P}_\partial \mathbf{q} \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &= \langle \mathbf{P}_\partial u - u, \mathbf{q} \cdot \mathbf{n} - \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &= -T_4, \end{aligned}$$

Finally, let us show that $T_5 = 0$. By the definition of the numerical trace $\widehat{\mathbf{q}}_h$, (1.4b),

$$\begin{aligned} T_5 &= \sum_{K \in \Omega_h} \langle \mathbb{P}u - u, (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial K} \\ &= \sum_{K \in \Omega_h} \langle \mathbb{P}u - u, (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{e_K^\tau} \\ &= 0, \end{aligned}$$

by the definition of the projection \mathbb{P} , (2.8b). This completes the proof. \square

3.3.2. *Proof of Theorem 2.6: Superconvergence of u_h .* Since

$$\|u - u_h\|_{H^{-s}(\Omega_h)} = \sup_{\theta \in \mathcal{C}_0^\infty(\Omega)} \frac{(\mathbb{P}u - u_h, \theta)_\Omega}{\|\theta\|_{H^s(\Omega)}}$$

we need to estimate the number $(\mathbb{P}u - u_h, \theta)_\Omega$. It is expressed in a suitable way in the following auxiliary result. Let us recall that \mathbf{P}^{k-1} is defined by (2.10) for $k \geq 1$, and is $\mathbf{P}^{k-1} \equiv 0$ for $k = 0$.

Lemma 3.8. *We have*

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} + (\mathbf{P}^{k-1}\nabla\varphi - \nabla\varphi, \mathbf{q} - \mathbf{\Pi}\mathbf{q})_{\Omega_h} \\ &\quad - \sum_{K \in \Omega_h} \tau^{-1} \langle \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} - \mathbf{P}_\partial \mathbf{q} \cdot \mathbf{n}, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} - \mathbf{P}_\partial \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{e_K^\tau}. \end{aligned}$$

Assume that $k \geq 1$. Then, applying the Cauchy-Schwarz inequality and using the estimate of $\mathbf{q} - \mathbf{q}_h$ in Theorem 2.5, and the approximation properties of the

projections \mathbb{P}^{k-1} and \mathbb{P} , (v) in Proposition 2.1 and (β) in Corollary 3.6, we readily obtain

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &\leq C h^{r+1} |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} h^{s+1} |\boldsymbol{\psi}|_{\mathbf{H}^{s+1}(\Omega_h)} \\ &\quad + C h^{r+1} |\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} h^{s+1} |\nabla \varphi|_{\mathbf{H}^{s+1}(\Omega_h)} \\ &\quad + \max_{K \in \Omega_h} \frac{1}{h_K \tau_K} C h^{r+1} |f|_{H^r(\Omega_h)} h^{s+1} |\theta|_{H^s(\Omega_h)}, \end{aligned}$$

where $r, s \in [0, k-1]$. Since $\kappa := \max_{K \in \Omega_h} \frac{1}{h_K \tau_K}$ and using the elliptic regularity assumption (2.14), we get

$$(\mathbb{P}u - u_h, \theta)_{\Omega_h} \leq C \mathfrak{C}_\kappa^{r,s}(\mathbf{q}) h^{r+s+2} |\theta|_{H^s(\Omega_h)}.$$

This completes the proof of Theorem 2.6 for $k \geq 1$.

In the case $k = 0$, we have that

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &\leq C h |\mathbf{q}|_{\mathbf{H}^1(\Omega_h)} h |\boldsymbol{\psi}|_{\mathbf{H}^1(\Omega_h)} \\ &\quad + C h |\mathbf{q}|_{\mathbf{H}^1(\Omega_h)} |\nabla \varphi|_{\mathbf{H}^1(\Omega_h)} \\ &\quad + \kappa C h |f|_{L^2(\Omega_h)} h |\theta|_{L^2(\Omega_h)}, \end{aligned}$$

and, after using the elliptic regularity assumption (2.14), we get

$$(\mathbb{P}u - u_h, \theta)_{\Omega_h} \leq C \mathcal{C}_\kappa(\mathbf{q}) h |\theta|_{L^2(\Omega_h)}.$$

Finally let us consider the case $k = 0$ and $f = 0$. By the identity (v) of Proposition 2.1 we have that $\mathbf{\Pi}\boldsymbol{\sigma} \cdot \mathbf{n} = \mathbb{P}_\partial \boldsymbol{\sigma} \cdot \mathbf{n}$, and by the identity (vi) of Proposition 2.1 we have that $\mathbf{\Pi}\mathbf{q} = \mathbf{\Pi}^{\text{RT}} \mathbf{q}$. This implies that

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} + \langle \varphi, (\mathbf{q} - \mathbf{\Pi}\mathbf{q}) \cdot \mathbf{n} \rangle_{\partial\Omega} \\ &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}^{\text{RT}} \boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} + \langle \varphi - \mathbb{P}_\partial \varphi, \mathbf{q}_\mathbf{N} - \mathbb{P}_\partial \mathbf{q}_\mathbf{N} \rangle_{\partial\Omega_N} \end{aligned}$$

by the adjoint equation (2.13c) and the boundary condition (1.1d). As a consequence, we get

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &\leq C h |\mathbf{q}|_{\mathbf{H}^1(\Omega_h)} h |\boldsymbol{\psi}|_{\mathbf{H}^1(\Omega_h)} \\ &\quad + C h |\mathbf{q}_\mathbf{N}|_{H^1(\partial\Omega_N)} h |\varphi|_{H^1(\partial\Omega_N)}, \end{aligned}$$

and since

$$|\varphi|_{H^1(\partial\Omega_N)} \leq C |\varphi|_{H^2(\Omega)},$$

by the elliptic regularity assumption (2.14), we get

$$(\mathbb{P}u - u_h, \theta)_{\Omega_h} \leq C \mathfrak{C}_\kappa^{0,0}(\mathbf{q}) h^2 |\theta|_{L^2(\Omega)}.$$

It remains to prove Lemma 3.8.

Proof. By the adjoint equation (2.13b), we have that

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbb{P}u - u_h, \nabla \cdot \boldsymbol{\psi})_{\Omega_h} \\ &= (u - u_h, \nabla \cdot \mathbf{\Pi}\boldsymbol{\psi})_{\Omega_h} - \langle \mathbb{P}_\partial u - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \end{aligned}$$

by the identity (α) of Corollary 3.6 with $(\boldsymbol{\sigma}, \zeta) := (\boldsymbol{\psi}, u - u_h)$. By the error equation (3.17a) with $\mathbf{v} := \mathbf{\Pi}\boldsymbol{\psi}$, we get

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi})_{\Omega_h} + \langle u - \widehat{u}_h, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad - \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} + (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \boldsymbol{\psi})_{\Omega_h} \\ &\quad + \langle u - \widehat{u}_h, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h} - \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \end{aligned}$$

and, by the adjoint equation (2.13a),

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} - (\mathbf{q} - \mathbf{q}_h, \nabla\varphi)_{\Omega_h} \\ &\quad + \langle u - \widehat{u}_h, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h} - \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}. \end{aligned}$$

By the orthogonality property (ii) of Proposition 2.1, we get that

$$\begin{aligned} (\mathbb{P}u - u_h, \theta)_{\Omega_h} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi})_{\Omega_h} + (\mathbf{q} - \mathbf{\Pi}\mathbf{q}, \mathbb{P}^{k-1}\nabla\varphi - \nabla\varphi)_{\Omega_h} \\ &\quad - (\mathbf{\Pi}\mathbf{q} - \mathbf{q}_h, \nabla\varphi)_{\Omega_h} + \langle u - \widehat{u}_h, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &\quad - \langle \mathbb{P}_{\partial}u - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}. \end{aligned}$$

If we denote by T the last three terms of the above right-hand side, we see that, after some simple algebraic manipulations, we can write $T = \sum_{i=1}^4 T_i$, where

$$\begin{aligned} T_1 &= - \langle \widehat{u}_h - u_h, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_2 &= - \langle \mathbb{P}_{\partial}u - u, (\mathbf{\Pi}\boldsymbol{\psi} - \boldsymbol{\psi}) \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_3 &= - \langle \widehat{u}_h - u, \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h}, \\ T_4 &= - (\mathbf{\Pi}\mathbf{q} - \mathbf{q}_h, \nabla\varphi)_{\Omega_h}. \end{aligned}$$

By the definition of the numerical trace \widehat{u}_h , (1.4a) and (1.6), we have that

$$\begin{aligned} T_1 &= - \sum_{K \in \Omega_h} \langle \widehat{u}_h - u_h, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} - \mathbb{P}_{\partial}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{e_K^{\tau}} \\ &= - \sum_{K \in \Omega_h} \tau^{-1} \langle \mathbf{\Pi}\mathbf{q} \cdot \mathbf{n} - \mathbb{P}_{\partial}\mathbf{q} \cdot \mathbf{n}, \mathbf{\Pi}\boldsymbol{\psi} \cdot \mathbf{n} - \mathbb{P}_{\partial}\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{e_K^{\tau}}, \end{aligned}$$

by Proposition 2.2.

It remains to show that $T_2 + T_3 + T_4 = 0$. By the definition of the projection \mathbb{P}_{∂} , (1.5),

$$\begin{aligned} T_2 &= - \langle \mathbb{P}_{\partial}u - u, -\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_h} \\ &= - \langle \mathbb{P}_{\partial}u - u, -\boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_D} \\ &= \langle \mathbb{P}_{\partial}u - u, \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_D}. \end{aligned}$$

By the definition of the numerical trace \widehat{u}_h , (1.4a),

$$\begin{aligned} T_3 &= - \langle \widehat{u}_h - u, \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega} \\ &= - \langle \widehat{u}_h - u, \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_D} \\ &= - \langle \mathbb{P}_{\partial}u - u, \boldsymbol{\psi} \cdot \mathbf{n} \rangle_{\partial\Omega_D} \\ &= - T_2. \end{aligned}$$

Next, we show that $T_4 = 0$. Integrating by parts, we obtain

$$\begin{aligned} T_4 &= (\nabla \cdot (\mathbf{\Pi} \mathbf{q} - \mathbf{q}_h), \varphi)_{\Omega_h} - \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}_h) \cdot \mathbf{n}, \varphi \rangle_{\partial \Omega_h} \\ &= (\nabla \cdot (\mathbf{q} - \mathbf{q}_h), \mathbb{P} \varphi)_{\Omega_h} + \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}) \cdot \mathbf{n}, \mathbb{P} \varphi \rangle_{\partial \Omega_h} \\ &\quad - \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}_h) \cdot \mathbf{n}, \varphi \rangle_{\partial \Omega_h} \end{aligned}$$

by the identity (α) of Corollary 3.6 with $(\boldsymbol{\sigma}, \zeta) := (\mathbf{q} - \mathbf{q}_h, \varphi)$. By the error equation (3.17b) with $\omega := \mathbb{P} \varphi$,

$$\begin{aligned} T_4 &= \langle \mathbb{P} \varphi, (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial \Omega_h} + \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}) \cdot \mathbf{n}, \mathbb{P} \varphi \rangle_{\partial \Omega_h} \\ &\quad - \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}_h) \cdot \mathbf{n}, \varphi \rangle_{\partial \Omega_h} \\ &= \langle \mathbb{P} \varphi, (\widehat{\mathbf{q}}_h - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial \Omega_h} + \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}) \cdot \mathbf{n}, \mathbb{P} \varphi \rangle_{\partial \Omega_h} \\ &\quad - \langle (\mathbf{\Pi} \mathbf{q} - \mathbf{q}_h) \cdot \mathbf{n}, \mathbb{P} \varphi \rangle_{\partial \Omega_h}, \end{aligned}$$

by the definition of the projection \mathbb{P}_∂ , (1.5), the definition of the projection \mathbb{P} , (2.8b), and the definition of the numerical trace $\widehat{\mathbf{q}}_h$, (1.4b). Hence

$$T_4 = \langle \mathbb{P} \varphi, (\widehat{\mathbf{q}}_h - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial \Omega_h} = \langle \mathbb{P} \varphi, (\widehat{\mathbf{q}}_h - \mathbf{q}) \cdot \mathbf{n} \rangle_{\partial \Omega} = 0,$$

by the adjoint equation (2.13c) and the equation (1.3c) for $\widehat{\mathbf{q}}_h$.

This completes the proof. \square

3.3.3. Proof of Theorem 2.8: Superconvergence of \widehat{u}_h . To prove this theorem, let us begin by estimating $\|\mathbb{P} \partial u - \widehat{u}_h\|_{L^2(e)}$ for each face e of each simplex K . For the face e_K^τ , we have that, by definition of the projection \mathbb{P} , (2.8),

$$\begin{aligned} \|\mathbb{P} \partial u - \widehat{u}_h\|_{L^2(e_K^\tau)} &= \|\mathbb{P} u - \widehat{u}_h\|_{L^2(e_K^\tau)} \\ &\leq \|\mathbb{P} u - u_h\|_{L^2(e_K^\tau)} + \|u_h - \widehat{u}_h\|_{L^2(e_K^\tau)} \\ &\leq \|\mathbb{P} u - u_h\|_{L^2(e_K^\tau)} + C \frac{h_K^{r+1/2}}{\tau_K} |f|_{H^r(K)} \end{aligned}$$

by Proposition 2.2 and the identity (v) of Proposition 2.1. By using a classical inverse inequality, we can conclude that

$$h_K^{1/2} \|\mathbb{P} \partial u - \widehat{u}_h\|_{L^2(e_K^\tau)} \leq C \left(\|\mathbb{P} u - u_h\|_{L^2(K)} + \frac{h_K^{r+1}}{\tau_K} |f|_{H^r(K)} \right).$$

Now we consider the error in the faces e of K which are different from the face e_K^τ . By the error equation (3.17a), we have that, for all $\mathbf{v} \in \mathcal{P}^k(K)$,

$$\begin{aligned} \langle \widehat{u}_h - \mathbb{P} \partial u, \mathbf{v} \cdot \mathbf{n} \rangle_{\partial K \setminus e_K^\tau} &= (\mathbf{c}(\mathbf{q} - \mathbf{q}_h), \mathbf{v})_K - (\mathbb{P} u - u_h, \nabla \cdot \mathbf{v})_K \\ &\quad - \langle \widehat{u}_h - \mathbb{P} \partial u, \mathbf{v} \cdot \mathbf{n} \rangle_{e_K^\tau}. \end{aligned}$$

Taking $\mathbf{v} := \mathbf{Z}$ given by Lemma 3.2 with $z = \widehat{u}_h - \mathbb{P} \partial u$, we obtain that

$$\begin{aligned} \|\widehat{u}_h - \mathbb{P} \partial u\|_{L^2(\partial K \setminus e_K^\tau)} &\leq C (h_K^{1/2} \|\mathbf{q} - \mathbf{q}_h\|_{L^2(K)} + h_K^{-1/2} \|\mathbb{P} u - u_h\|_{L^2(K)} \\ &\quad + \|\widehat{u}_h - \mathbb{P} \partial u\|_{L^2(e_K^\tau)}), \end{aligned}$$

and using the estimate for the error in e_K^τ ,

$$\begin{aligned} h_K^{1/2} \|\widehat{u}_h - \mathbb{P} \partial u\|_{L^2(\partial K \setminus e_K^\tau)} &\leq C (\|\mathbb{P} u - u_h\|_{L^2(K)} + h_K \|\mathbf{q} - \mathbf{q}_h\|_{L^2(K)}^2 \\ &\quad + \frac{h_K^{r+1}}{\tau_K} |f|_{H^r(K)}). \end{aligned}$$

As a consequence

$$\begin{aligned} \|\mathbf{P}_{\partial} u - \widehat{u}_h\|_{L^2(\mathcal{E}_h; h)} &\leq C \left(\|\mathbb{P}u - u_h\|_{L^2(\Omega_h)} + h \|\mathbf{q} - \mathbf{q}_h\|_{L^2(\Omega_h)} \right. \\ &\quad \left. + \kappa h^{r+2} |f|_{H^r(\Omega_h)} \right), \end{aligned}$$

where $\kappa := \max_{K \in \Omega_h} 1/(\tau_K h_K)$. The result now follows from Theorems 2.6, 2.5 and 2.4 (i). This completes the proof of Theorem 2.8.

3.3.4. *Proof of Theorem 2.9: The error estimate for u_h^* .* By the definition of u_h^* , (2.15a), we have that

$$\|u - u_h^*\|_{L^2(K)} \leq \|\bar{u} - \bar{u}_h\|_{L^2(K)} + \|\tilde{u} - \tilde{u}_h\|_{L^2(K)},$$

where \bar{u} is defined in (2.15b) and $\tilde{u} = u - \bar{u}$. We estimate each of the two terms of the right-hand side separately.

We begin by estimating the second term. Since, by Poincaré's inequality, we have

$$\|\tilde{u} - \tilde{u}_h\|_{L^2(K)} \leq C h_K \|\nabla(\tilde{u} - \tilde{u}_h)\|_{L^2(K)},$$

it is enough to estimate the error in the gradient. To do that, we note that, by the definition of \tilde{u}_h , (2.15c), we have

$$(\mathbf{a} \nabla(\tilde{u} - \tilde{u}_h), \nabla w)_K = -\langle w, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial K} \quad \forall w \in \mathcal{P}_0^{k+1}(K).$$

Then

$$\begin{aligned} \|\nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h)\|_{L^2(K; \mathbf{a})}^2 &= (\mathbf{a} \nabla(\tilde{u} - \tilde{u}_h), \nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h))_K \\ &\quad + (\mathbf{a} \nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}), \nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h))_K \\ &= -\langle \mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial K} \\ &\quad + (\mathbf{a} \nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}), \nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h))_K. \end{aligned}$$

Let us estimate the first term of the right-hand side. For any arbitrary $\omega \in \mathcal{P}_0^{k+1}(K)$, we have

$$\begin{aligned} \langle \omega, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial K} &= \langle \omega, (\mathbf{q} - \mathbf{q}_h) \cdot \mathbf{n} \rangle_{\partial K} + \langle \omega, (\mathbf{q}_h - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial K} \\ &= \sum_{i=1}^3 T_i, \end{aligned}$$

where

$$\begin{aligned} T_1 &= (\nabla \omega, \mathbf{q} - \mathbf{q}_h)_K, \\ T_2 &= (\omega, \nabla \cdot (\mathbf{q} - \mathbf{q}_h))_K, \\ T_3 &= (\omega, (\mathbf{q}_h - \widehat{\mathbf{q}}_h) \cdot \mathbf{n})_{\partial K}. \end{aligned}$$

By using Cauchy-Schwarz inequality, we get that

$$T_1 \leq \|\nabla \omega\|_{L^2(K; \mathbf{a})} \|\mathbf{q} - \mathbf{q}_h\|_{L^2(K; \mathbf{c})}.$$

By using the definition of the Raviart-Thomas projection $\mathbf{\Pi}^{\text{RT}}$, (2.9), and by using its commutativity property, we get that, for any $r \in [0, k]$,

$$\begin{aligned} T_2 &= (\omega, f - \mathbf{P}f)_K + (\omega, \nabla \cdot (\mathbf{\Pi}^{\text{RT}} \mathbf{q} - \mathbf{q}_h))_K \\ &\leq \|\omega\|_{L^2(K)} (h_K^r |f|_{H^r(K)} + h_K^{-1} \|\mathbf{q} - \mathbf{q}_h\|_{L^2(K; \mathbf{c})}) \\ &\leq \|\nabla \omega\|_{L^2(K; \mathbf{a})} (h_K^{r+1} |f|_{H^r(K)} + \|\mathbf{q} - \mathbf{q}_h\|_{L^2(K; \mathbf{c})}) \end{aligned}$$

by Poincaré's inequality. Finally, by the definition of the numerical trace $\widehat{\mathbf{q}}_h$, (1.4b),

$$\begin{aligned} T_3 &= \|\omega\|_{L^2(e_K^r)} \|\mathbf{q}_h - \widehat{\mathbf{q}}_h\|_{L^2(e_K^r)} \cdot \mathbf{n} \\ &\leq C \|\omega\|_{L^2(e_K^r)} h_K^{r+1/2} |f|_{H^r(K)}, \end{aligned}$$

by Proposition 2.2 and identity (v) of Proposition 2.1. Applying a simple inverse inequality, we get

$$T_3 \leq C \|\nabla\omega\|_{L^2(K;\mathbf{a})} h_K^{r+1} |f|_{H^r(K)}.$$

As a consequence,

$$\langle \omega, (\mathbf{q} - \widehat{\mathbf{q}}_h) \cdot \mathbf{n} \rangle_{\partial K} \leq C \|\nabla\omega\|_{L^2(K;\mathbf{a})} (\|\mathbf{q} - \mathbf{q}_h\|_{L^2(K;\mathbf{c})} + h_K^{r+1} |f|_{H^r(K)}).$$

This implies that

$$\begin{aligned} \|\nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h)\|_{L^2(K;\mathbf{a})} &\leq \|\nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u})\|_{L^2(K)} \\ &\quad + C (\|\mathbf{q} - \mathbf{q}_h\|_{L^2(K;\mathbf{c})} + h_K^{r+1} |f|_{H^r(K)}), \end{aligned}$$

and so,

$$\begin{aligned} \|\mathbf{P}^{k+1}\tilde{u} - \tilde{u}_h\|_{L^2(\Omega_h;\mathbf{a})} &\leq Ch \|\nabla(\mathbf{P}^{k+1}\tilde{u} - \tilde{u})\|_{L^2(\Omega_h)} \\ &\quad + C (h \|\mathbf{q} - \mathbf{q}_h\|_{L^2(\Omega_h;\mathbf{c})} + h^{r+2} |f|_{H^r(\Omega_h)}) \\ &\leq Ch^{r+2} (|\mathbf{q}|_{\mathbf{H}^{r+1}(\Omega_h)} + |u|_{H^{r+2}(\Omega_h)}), \end{aligned}$$

by Theorem 2.5 and the well-known approximation properties of \mathbf{P}^{k+1} .

Let us now estimate the error $\bar{u} - \bar{u}_h$. We begin by considering the case $k \geq 1$. In this case, since $\bar{u} - \bar{u}_h = \overline{\mathbb{P}(u - u_h)}$, we get

$$\|\bar{u} - \bar{u}_h\|_{L^2(K)} \leq \|\mathbb{P}u - u_h\|_{L^2(K)} \leq C \mathfrak{C}_\kappa^{r,0}(\mathbf{q}) h^{r+2}$$

by Theorem 2.6. Note that by Theorem 2.4, $\mathbf{P}^{k-1}u_h$ is independent of the value of the local stabilization parameters τ . This implies that the same is true for \bar{u}_h and so, we get that

$$\|\bar{u} - \bar{u}_h\|_{L^2(K)} \leq C \mathfrak{C}_0^{r,0}(\mathbf{q}) h^{r+2}.$$

It remains to consider the case $k = 0$ and $f = 0$. We have that

$$\begin{aligned} \bar{u} - \bar{u}_h &= \frac{1}{|K|} \int_K u(x) dx - \frac{1}{d} \sum_{e \in \partial K} \widehat{u}_h|_e \\ &= \frac{1}{|K|} \int_K (u - \mathbf{P}^1 u)(x) dx \\ &\quad + \frac{1}{|K|} \int_K \mathbf{P}^1 u(x) dx - \frac{1}{d} \sum_{e \in \partial K} \frac{1}{|e|} \int_e \mathbf{P}_\partial(\mathbf{P}^1 u) d\gamma \\ &\quad - \frac{1}{d} \sum_{e \in \partial K} \frac{1}{|e|} \int_e (\mathbf{P}_\partial(u - \mathbf{P}^1 u)) d\gamma \\ &\quad - \frac{1}{d} \sum_{e \in \partial K} \frac{1}{|e|} \int_e (\widehat{u}_h - \mathbf{P}_\partial u) d\gamma. \end{aligned}$$

Since, for any function $\omega \in \mathcal{P}^1(K)$, we have that

$$\frac{1}{|K|} \int_K \omega(x) dx - \frac{1}{d} \sum_{e \in \partial K} \frac{1}{|e|} \int_e \mathbf{P}_\partial \omega d\gamma = 0,$$

we readily obtain that

$$\begin{aligned} \|\bar{u} - \bar{u}_h\|_{L^2(K)} &\leq C \|u - \mathbf{P}^1 u\|_{L^2(K)} + C h_K |u - \mathbf{P}^1 u|_{H^1(K)} \\ &\quad + C h_k \|\hat{u}_h - \mathbf{P}_{\partial} u\|_{L^2(\partial K)} \end{aligned}$$

and so,

$$\|\bar{u} - \bar{u}_h\|_{L^2(\Omega_h)} \leq C h^2 (\mathcal{C}_\kappa(\mathbf{q}) + |u|_{H^2(\Omega_h)}).$$

Since, by Theorem 2.4, λ_h is independent of the value of the local stabilization parameter τ , so is \bar{u}_h and so

$$\|\bar{u} - \bar{u}_h\|_{L^2(K)} \leq C h^2 (\mathcal{C}_0(\mathbf{q}) + |u|_{H^2(\Omega_h)}).$$

This completes the proof of Theorem 2.9.

4. NUMERICAL EXPERIMENTS

In this section, we carry out numerical experiments to validate the theoretical convergence properties of the SCDG_k method.

To do that, we use uniform meshes obtained by discretizing $\Omega = (-\frac{1}{2}, \frac{1}{2}) \times (-\frac{1}{2}, \frac{1}{2})$ with squares of side 2^{-l} which are then divided into two triangles as indicated in Fig. 1; the resulting mesh is denoted by “mesh= l ”.

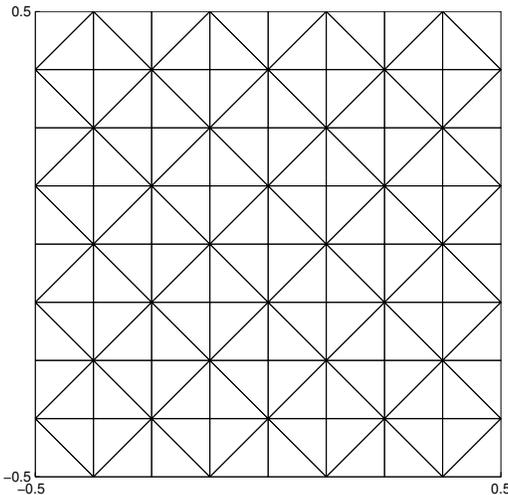


FIGURE 1. Example of a mesh with $h = 1/2^3$.

The test problem is obtained by taking $\partial\Omega_N = \emptyset$, $\mathbf{c} = \mathbf{I}$ and choosing g and f so that the exact solution is $u(x, y) = \cos(\pi x) \cos(\pi y)$ on the domain Ω . The history of convergence of the SCDG method with

$$\tau_K = 1/h = 2^l,$$

on the “mesh= l ”, is displayed in Table 5 for polynomials of degree $k = 0$, $k = 1$ and $k = 2$. We observe optimal convergence rates of the quantities $\|u - u_h\|_{L^2(\Omega)}$

and $\|q - q_h\|_{L^2(\Omega)}$ for $k = 0, 1, 2$ as predicted by Theorems 2.5 and 2.7. We also see that $\|P_{\partial}u - \lambda_h\|_{L^2(\mathcal{E}_h;h)}$ and $\|u - u_h^*\|_{L^2(\Omega)}$ superconverges with rate $O(h^{k+2})$ for $k = 1, 2$ just as predicted by Theorems 2.8 and 2.9. These results do not guarantee that these quantities are superconvergent if $k = 0$ and $f \neq 0$. Since we do not observe superconvergence, we can conclude that the theoretical results for such a case are actually sharp.

TABLE 5. History of convergence of the SCDG_k method.

k	mesh ℓ	$\ u - u_h\ _{L^2(\Omega)}$		$\ q - q_h\ _{L^2(\Omega)}$		$\ P_{\partial}u - \lambda_h\ _{L^2(\mathcal{E}_h;h)}$		$\ u - u_h^*\ _{L^2(\Omega)}$	
		error	order	error	order	error	order	error	order
0	1	.11e+1	-	.17e+1	-	.28e-0	-	.22e-0	-
	2	.36e-0	1.54	.78e-0	1.12	.92e-1	1.61	.57e-1	1.96
	3	.12e-0	1.50	.41e-0	0.94	.35e-1	1.37	.19e-1	1.57
	4	.53e-1	1.29	.21e-0	0.97	.14e-1	1.21	.79e-2	1.27
	5	.24e-1	1.13	.10e-0	0.98	.69e-2	1.15	.36e-2	1.14
	6	.12e-1	1.05	.53e-1	0.99	.32e-2	1.12	.17e-2	1.10
1	1	.21e-0	-	.23e-0	-	.31e-1	-	.21e-1	-
	2	.43e-1	2.27	.12e-0	0.94	.75e-2	2.02	.40e-2	2.38
	3	.78e-2	2.47	.31e-1	1.94	.10e-2	2.96	.53e-3	2.91
	4	.17e-2	2.19	.79e-2	1.99	.12e-3	3.00	.68e-4	2.98
	5	.42e-3	2.05	.20e-2	2.00	.15e-4	3.01	.85e-5	2.99
	6	.10e-3	2.01	.50e-3	2.00	.19e-5	3.00	.11e-6	2.99
2	1	.68e-1	-	.89e-1	-	.72e-2	-	.81e-2	-
	2	.38e-2	4.12	.91e-2	3.29	.41e-3	4.12	.40e-3	4.35
	3	.32e-3	3.58	.12e-2	2.96	.27e-4	3.93	.25e-4	3.97
	4	.32e-4	3.31	.15e-3	2.98	.18e-5	3.96	.16e-5	3.99
	5	.37e-5	3.12	.19e-4	2.99	.11e-6	3.98	.10e-6	4.00
	6	.45e-7	3.01	.23e-5	3.00	.70e-8	3.99	.63e-8	4.00

Next we explore the effect of the size of τ_K on the quality of the approximation. In table 6, we see that as τ diminishes the quality of the approximation to u deteriorates. However, the effect of taking $\tau = 1/h^2$ or $\tau = 1/h$ is almost negligible especially when the grids are not coarse. We also see that the order of convergence is $k + 1$ for $\tau = 1/h^2, \tau = 1/h$ and $\tau = 1$, but it is only k for $\tau = h$. This is in perfect agreement with Corollary 2.7.

We end with an example where the exact solution is harmonic, that is, $\mathbf{c} = \mathbf{I}$, $f \equiv 0$, and display the convergence rates for $k = 0$ in Table 7. We take $\partial\Omega_N = \emptyset$ and choose g and so that $u(x, y) = e^x \sin(y)$ is the solution. We see that the quantities $\|P_{\partial}u - \lambda_h\|_{L^2(\mathcal{E}_h;h)}$ and $\|u - u_h^*\|_{L^2(\Omega)}$ superconverge with rate $O(h^2)$ as our theoretical results predict.

5. CONCLUDING REMARKS

The error analysis carried out here for the SCDG_k method also holds for the hybridized versions of the RT_K and the BDM_k methods. We simply have to replace the local space $\mathcal{P}^k(K) \times \mathcal{P}^k(K)$ by the local space $\mathbf{V}(K) \times W(K)$ given by Table 1, use the definition of the local stabilization parameter τ given in Table 2, and suitably define the projection $(\mathbf{\Pi}, \mathbb{P})$. Indeed, with such changes, the first four properties of Proposition 2.1, on which the whole analysis is based upon, do hold. For this reason, we can consider this analysis to be a *unifying* analysis of these three methods.

TABLE 6. Effect of τ on the convergence of $\|u - u_h\|_{L^2(\Omega)}$.

k	mesh	$\tau = 1/h^2$		$\tau = 1/h$		$\tau = 1$		$\tau = h$	
	ℓ	error	order	error	order	error	order	error	order
0	1	.61e+0	-	.11e+1	-	.21e+1	-	.40e+1	-
	2	.21e+0	1.50	.36e-0	1.54	.11e+1	0.88	.42e+1	-0.06
	3	.95e-1	1.17	.12e-0	1.50	.57e-0	0.97	.43e+1	-0.02
	4	.46e-1	1.04	.53e-1	1.29	.28e-0	1.00	.43e+1	-0.00
	5	.23e-1	1.02	.24e-1	1.13	.14e-0	1.01	.43e+1	0.00
	6	.11e-1	1.01	.12e-1	1.05	.70e-1	1.01	.43e+1	0.00
1	1	.15e+0	-	.21e-0	-	.35e-0	-	.67e-0	-
	2	.27e-1	2.47	.43e-1	2.27	.14e-0	1.34	.54e-0	0.29
	3	.66e-2	2.06	.78e-2	2.47	.35e-1	1.98	.28e-0	0.97
	4	.16e-2	2.00	.17e-2	2.19	.88e-2	2.00	.14e-0	0.99
	5	.41e-3	2.00	.42e-3	2.05	.22e-2	2.00	.67e-1	1.00
	6	.10e-3	2.00	.10e-3	2.01	.60e-3	2.00	.35e-1	1.00
2	1	.33e-1	-	.68e-1	-	.14e-0	-	.28e-0	-
	2	.19e-2	4.10	.38e-2	4.12	.14e-1	3.31	.56e-1	2.33
	3	.23e-3	3.07	.32e-3	3.58	.18e-2	2.96	.14e-1	1.98
	4	.28e-4	3.00	.32e-4	3.31	.23e-3	2.99	.35e-2	2.00
	5	.36e-5	3.00	.37e-5	3.12	.28e-4	3.00	.88e-3	2.00
	6	.45e-6	3.00	.45e-7	3.01	.35e-5	3.00	.22e-3	2.00

TABLE 7. History of convergence for a harmonic exact solution.

k	mesh	$\ u - u_h\ _{L^2(\Omega)}$		$\ q - q_h\ _{L^2(\Omega)}$		$\ P_{\partial} u - \lambda_h\ _{L^2(\varepsilon_h; h)}$		$\ u - u_h^*\ _{L^2(\Omega)}$	
	ℓ	error	order	error	order	error	order	error	order
0	1	.17e-0	-	.22e-0	-	.29e-1	-	.23e-1	-
	2	.87e-1	0.94	.11e-0	0.96	.79e-2	1.84	.62e-2	1.87
	3	.44e-1	0.99	.57e-1	0.98	.21e-2	1.90	.16e-2	1.93
	4	.22e-1	0.99	.29e-1	0.99	.55e-3	1.96	.41e-3	1.97
	5	.11e-1	1.00	.14e-1	1.00	.10e-3	1.98	.10e-3	1.99
	6	.55e-2	1.00	.72e-2	1.00	.35e-4	2.00	.26e-4	2.00

A study of the optimal way to choose the local stabilization parameter τ falls beyond the scope of this paper and will be carried out elsewhere. Extensions of these results to more general second-order elliptic equations and other boundary conditions are straightforward. The extension of these results to the case of hanging nodes, variable-degree approximations and curved domains constitute the subject of ongoing work.

REFERENCES

- [1] D. N. Arnold and F. Brezzi, *Mixed and nonconforming finite element methods: implementation, postprocessing and error estimates*, RAIRO Modél. Math. Anal. Numér. **19** (1985), 7–32.
- [2] D. N. Arnold, F. Brezzi, B. Cockburn, and L. D. Marini, *Unified analysis of discontinuous Galerkin methods for elliptic problems*, SIAM J. Numer. Anal. **39** (2002), 1749–1779.
- [3] J. H. Bramble and J. Xu, *A local post-processing technique for improving the accuracy in mixed finite-element approximations*, SIAM J. Numer. Anal. **26** (1989), no. 6, 1267–1275.
- [4] F. Brezzi, J. Douglas, Jr., and L. D. Marini, *Two families of mixed finite elements for second order elliptic problems*, Numer. Math. **47** (1985), 217–235.

- [5] P. Castillo, B. Cockburn, I. Perugia, and D. Schötzau, *An a priori error analysis of the local discontinuous Galerkin method for elliptic problems*, SIAM J. Numer. Anal. **38** (2000), 1676–1706.
- [6] F. Celiker and B. Cockburn, *Superconvergence of the numerical traces of discontinuous Galerkin and hybridized mixed methods for convection-diffusion problems in one space dimension*, Math. Comp. **76** (2007), 67–96.
- [7] B. Cockburn and B. Dong, *An analysis of the minimal dissipation local discontinuous Galerkin method for convection-diffusion problems*, J. Sci. Comput., To appear. Available as the I.M.A. preprint 2146, November 2006.
- [8] B. Cockburn and J. Gopalakrishnan, *A characterization of hybridized mixed methods for second order elliptic problems*, SIAM J. Numer. Anal. **42** (2004), 283–301.
- [9] B. Cockburn, J. Gopalakrishnan, and R. Lazarov, *Unified hybridization of discontinuous Galerkin, mixed and continuous Galerkin methods for second order elliptic problems*, Math. Comp., Submitted.
- [10] B. Cockburn and C.-W. Shu, *The local discontinuous Galerkin method for time-dependent convection-diffusion systems*, SIAM J. Numer. Anal. **35** (1998), 2440–2463.
- [11] J. Douglas, Jr. and J.E. Roberts, *Global estimates for mixed methods for second order elliptic equations*, Math. Comp. **44** (1985), 39–52.
- [12] Lucia Gastaldi and Ricardo H. Nochetto, *Sharp maximum norm error estimates for general mixed finite element approximations to second order elliptic equations*, RAIRO Modél. Math. Anal. Numér. **23** (1989), no. 1, 103–128. MR MR1015921 (91b:65125)
- [13] P. A. Raviart and J. M. Thomas, *A mixed finite element method for second order elliptic problems*, Mathematical Aspects of Finite Element Method (I. Galligani and E. Magenes, eds.), Lecture Notes in Math. 606, Springer-Verlag, New York, 1977, pp. 292–315.
- [14] Rolf Stenberg, *A family of mixed finite elements for the elasticity problem*, Numer. Math **53** (1988), 513–538.
- [15] Rolf Stenberg, *Postprocessing schemes for some mixed finite elements*, RAIRO Modél. Math. Anal. Numér. **25** (1991), 151–167.

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