Efficient computing in evolution equations with memory

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Soft tissue biomechanics:

- Complex biomaterials with non-trivial properties
- Non-linear elastic response
- Viscoelastic creep and continuous relaxation time-scales

Fung, Yuan-cheng. *Biomechanics: mechanical properties of living tissues*. Springer Science & Business Media, 2013.

Modeling approach #1: Simple parametric integer-order models



Standard Linear Solid model (SLS):

$$\sigma(t) + \tau_{\sigma} \frac{d\sigma(t)}{dt} = E[\epsilon(t) + \tau_{\epsilon} \frac{d\epsilon(t)}{dt}]$$
 3 parameters

- Linear elastic response
- Simplest model that accounts for viscoelastic creep and relaxation
- Two discreet relaxation time-scales



Modeling approach #2: More complex parametric integer-order models

QLV Kelvin-Zener solids in series



- Linear elastic response
- Can capture more complex viscoelastic creep and relaxation behavior
- Multiple discreet relaxation time-scales



- Many parameters to be estimated from few data
- High parametric sensitivity
- The response behaves like a stochastic process and depends on patient-specific factors, age, etc.

Fung, Yuan-cheng. *Biomechanics: mechanical properties of living tissues*. Springer Science & Business Media, 2013.

Modeling approach #3: Even more complex parametric integer-order models

QLV Kelvin-Zener solids in series



- Linear elastic response
- Can capture even more complex viscoelastic creep and relaxation behavior
- Continuous spectrum of relaxation time-scales

... Guess how many parameters?

Modeling approach #4: Fractional-order models

Fractional-order Kelvin-Zener solid



- Linear elastic response
- Can capture even more complex viscoelastic creep and relaxation behavior
- Continuous spectrum of relaxation time-scales

$$\sigma(t) + \tau^{\alpha C}_{\sigma \ 0} D^{\alpha}_{t} \sigma(t) = E \left[\epsilon(t) + \tau^{\alpha C}_{\epsilon \ 0} D^{\alpha}_{t} \epsilon(t) \right]$$

4 parameters!

Doehring, Todd C., et al. "Fractional order viscoelasticity of the aortic valve cusp: an alternative to quasilinear viscoelasticity." Journal of biomechanical engineering 127.4 (2005): 700-708. Näsholm, Sven Peter, and Sverre Holm. "On a fractional Zener elastic wave equation." Fractional Calculus and Applied Analysis 16.1 (2013): 26-50.

Modeling approach #4: More complex fractional-order models

Fractional-order Kelvin-Zener solid



- Linear elastic response
- Can capture even more complex viscoelastic creep and relaxation behavior
- Continuous spectrum of relaxation time-scales

$$\sigma(t) + \tau_{\sigma \ 0}^{\alpha C} D_t^{\alpha} \sigma(t) = E \left[\epsilon(t) + \tau_{\epsilon \ 0}^{\alpha C} D_t^{\alpha} \epsilon(t) \right]$$

4 parameters!



... fewer parameters to be estimated, hence less parametric sensitivity.

... combine fractional solids in series and you can fit an elephant in the room (e.g. smooth muscle activation in cerebral auto regulation).

The role of the fractional order

 $\beta \cong 0.7$ (higher value)



The fractional order effectively controls the balance between the conservative (elastic energy storage) and dissipative (viscoelastic energy dissipation) behaviors.







 $\sim \sim$

Great, but how do we compute now?

Discretize the derivatives

Grunwald-Letnikov approximation

$${}_{0}^{C}D_{t}^{\alpha}f(t) = \lim_{\Delta t \to 0} \Delta t^{-\alpha} \sum_{k=0}^{\infty} GL_{k}^{\alpha}f(t-k\Delta t), \quad GL_{k}^{a} := \frac{k-\alpha-1}{k}GL_{k-1}^{\alpha}$$

$$p(x,t) = p_{ext} + \frac{1 + \tau_{\epsilon}^{\alpha} \Delta t^{-\alpha}}{1 + \tau_{\sigma}^{\alpha} \Delta t^{-\alpha}} p^{E}(x,t) + \frac{\Delta t^{-\alpha}}{1 + \tau_{\sigma}^{\alpha} \Delta t^{-\alpha}} \sum_{k=0}^{\infty} GL_{k}^{\alpha} \{\tau_{\epsilon}^{\alpha} p^{E}(t-k\Delta t) - \tau_{\sigma}^{\alpha} p(t-k\Delta t)\}$$

- Low accuracy
- Excessive memory requirements, especially for advection dominated problems due to the CFL condition

...why should we discretize in the first place?
These are linear time-fractional evolution equations with well behaved forcing terms...
...we can use our good old Laplace transforms and solve them analytically!

Computing convolution integrals

$$p(x,t) = p_{ext} + p^{E}(x,t) + p^{V}(x,t)$$

$$p^{E}(x,t) = \left(\frac{\tau_{\epsilon}}{\tau_{\sigma}}\right)^{\alpha} \beta(\sqrt{A} - \sqrt{A_{0}})$$

$$p^{V}(x,t) = \frac{1}{\tau_{\sigma}} \left(1 - \left[\frac{\tau_{\epsilon}}{\tau_{\sigma}}\right]^{\alpha}\right) \int_{0}^{t} E_{\alpha,0} \left(-\left[\frac{t-\gamma}{\tau_{\sigma}}\right]^{\alpha}\right) p^{E}(\gamma) d\gamma$$

Cost of a naive implementation:

$$O(N^2)$$
 operations
 $O(N)$ evaluations of the kernel (e.g. Mittal-Leffler function)
 $O(N)$ active memory

Should be more accurate than the GL discretization but it doesn't buy us much in terms of computational efficiency.

Fast convolutions for evolution equations with memory

Analytical solution of FPDEs using Laplace transforms gives rise to convolution integrals:

$$\int_0^t f(t-\tau) g(\tau) d\tau, \qquad 0 \le t \le T$$

The fast convolution method of [1] allows to compute such integrals with spectral accuracy and

- $O(N \log N)$ operations,
- $O(\log N)$ evaluations of the Laplace transform $F = \mathcal{L}f$, none of f, and
- $O(\log N)$ active memory.

	$O(N^2)$	for a naive
VS	O(N)	quadrature
	O(N)	implementation

Example: In fractional-order viscoelastic constitute laws we need to compute a memory term:

$$\int_{0}^{t} E_{\alpha,0} \left(-\left[\frac{t-\gamma}{\tau_{\sigma}}\right]^{\alpha} \right) p^{e}(\gamma) d\gamma$$

The fast convolution method bypasses costly evaluations of the Mittag-Leffler kernel: $E_{\alpha}(x) = \sum_{j=0}^{\infty} \frac{x^j}{\Gamma(1+\alpha j)}$ and only requires simple function evaluations in the Laplace domain: $F(s) = \frac{1}{1+s^{\alpha}}$

...and only requires simple function evaluations in the Laplace domain:

Implementation

$$\begin{aligned} \int_{a}^{b} f(t-\tau) g(\tau) \, d\tau &= \int_{a}^{b} \frac{1}{2\pi i} \int_{\Gamma} F(\lambda) \, e^{(t-\tau)\lambda} \, d\lambda \, g(\tau) \, d\tau \\ &= \frac{1}{2\pi i} \int_{\Gamma} F(\lambda) \, e^{(t-b)\lambda} \int_{a}^{b} e^{(b-\tau)\lambda} \, g(\tau) \, d\tau \, d\lambda \approx \int_{a}^{b} \sum_{j=-N}^{N} w_{j} \, F(\lambda_{j}) \, e^{(t-\tau)\lambda_{j}} \, g(\tau) \, d\tau \ = \sum_{j=-N}^{N} w_{j} \, F(\lambda_{j}) \, e^{(t-b)\lambda_{j}} \, y(b,a,\lambda_{j}) \end{aligned}$$

where the inner integral, henceforth denoted by $y(b, a, \lambda)$, is recognized as the solution at time b of the scalar linear initial value problem

$$y' = \lambda y + g$$
, $y(a) = 0$.

The 2N + 1 differential equations (2.5) with $\lambda = \lambda_j$ are solved approximately by replacing the function g with its piecewise linear approximation and then solving exactly. Setting $g_n = g(a + n\Delta t)$, we get approximations $y_n \approx y(a + n\Delta t)$ recursively via

$$y_{n+1} = e^{\Delta t\lambda} y_n + h \int_0^1 e^{(1-\theta)\Delta t\lambda} \left(\theta g_{n+1} + (1-\theta)g_n\right) d\theta$$

(2.7)
$$= y_n + \frac{e^{\Delta t\lambda} - 1}{\Delta t\lambda} \left(\Delta t\lambda y_n + \Delta tg_n + \Delta t \frac{g_{n+1} - g_n}{\Delta t\lambda}\right) - \Delta t \frac{g_{n+1} - g_n}{\Delta t\lambda}.$$

To estimate the error, note that in total we approximate

$$\int_{a}^{b} f(t-\tau) g(\tau) d\tau \approx \int_{a}^{b} \widetilde{f}(t-\tau) \widetilde{g}(\tau) d\tau,$$

where \tilde{f} is the quadrature approximation to f constructed in the previous subsection, whose error is well under control, and \tilde{g} is the piecewise linear interpolant of g. (Higher-order interpolants of g might also be used, but we have not implemented such an extension.)

Implementation

Laplace transform $\int_{0}^{t_{n}} \underbrace{f(t_{n} - \tau)}_{g(\tau)} d\tau = \int_{0}^{t_{n}} \frac{1}{2\pi i} \int_{\Gamma} \underbrace{e^{(t_{n} - \tau)\lambda} F(\lambda)}_{e^{(t_{n} - \tau)\lambda} F(\lambda)} d\lambda g(\tau) d\tau$

$$\int_{t^{-}}^{t^{+}} f(t_n - \tau)g(\tau) d\tau = \int_{t^{-}}^{t^{+}} \frac{1}{2\pi i} \int_{\Gamma} e^{(t_n - \tau)\lambda} F(\lambda) d\lambda g(\tau) d\tau$$

$$\approx \int_{t^{-}}^{t^{+}} \sum_{k=-K}^{K} w_k e^{(t_n - \tau)\lambda_k} F(\lambda_k) g(\tau) d\tau$$

$$= \sum_{k=-K}^{K} w_k F(\lambda_k) e^{(t_n - t^{+})\lambda_k} \int_{t^{-}}^{t^{+}} e^{(t^{+} - \tau)\lambda_k} g(\tau) d\tau$$

$$= \sum_{k=-K}^{K} w_k F(\lambda_k) e^{(t_n - t^{+})\lambda_k} y(t^{+}, t^{-}, \lambda_k),$$

weights w_k and quadrature nodes λ_k given by

$$w_k = \frac{\tau}{2\pi i} \gamma'(k\tau) , \quad \lambda_k = \gamma(k\tau)$$

where $y(t^+, t^-, \lambda_k)$ is the solution at t^+ to the linear inhomogeneous ODE $y' = \lambda_k y + g$, $y(t^-) = 0$, $-K \le k \le K$

<u>Key ingredients:</u>

- Complex plane quadrature (Talbot contours)
- Local reduction to ODEs
- Adaptivity and smart memory management

····· Forward Euler, ···► Runge-Kutta, BDF, etc

"Filling the mosaic"

Matlab demo

$$\int_{a}^{b} f(t-\tau) d\tau$$

$$f(t) = \frac{1}{\sqrt{\pi t}} \quad \text{for which} \quad F(s) = s^{-1/2}$$